

# Heavy Ions at LHC

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# Outlook

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- QGP
- Event characterisation
- Soft probes
  - Interferometry
  - Multiplicity, Transverse energy, Energy density
  - Flow and correlations
- Hard Probes
  - Quarkonia
  - Jet quenching
    - High pt suppression ( $h^-$ ,  $D0$ ,  $J/\psi$ ,  $\Upsilon$ ,  $Z$ , ...)
    - Reconstructed jets
- Summary

# Quantum ChromoDynamics (QCD)

$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_j \bar{\psi}_j (i\gamma^\mu D_\mu + m_j) \psi_j$$

$$\text{where } G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$$

$$\text{and } D_\mu = \partial_\mu + ig A_\mu^a$$

That's it!

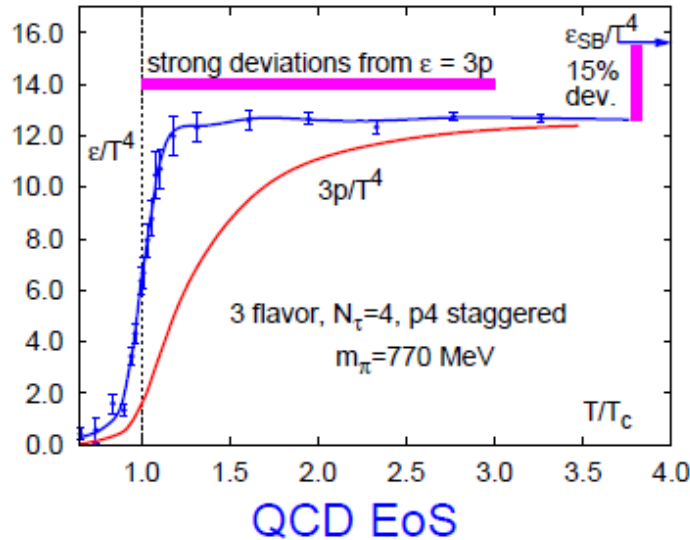
**QCD confinement – free quarks never observed !**

**QCD vacuum not well understood.**

**Heavy ions – study QCD at high temperature and density**

# Quark Gluon Plasma

- Lattice QCD: transition hadrons -> quarks and gluons



QGP is not ideal gas !

$$p = \frac{\varepsilon}{3} = \left( g_B + \frac{7}{8} g_F \right) \frac{\pi^2 T^4}{90}$$

$$g_B = 8_c * 2_s = 16$$

$$g_F = 3_f * 3_c * 2_s * 2_a = 36$$

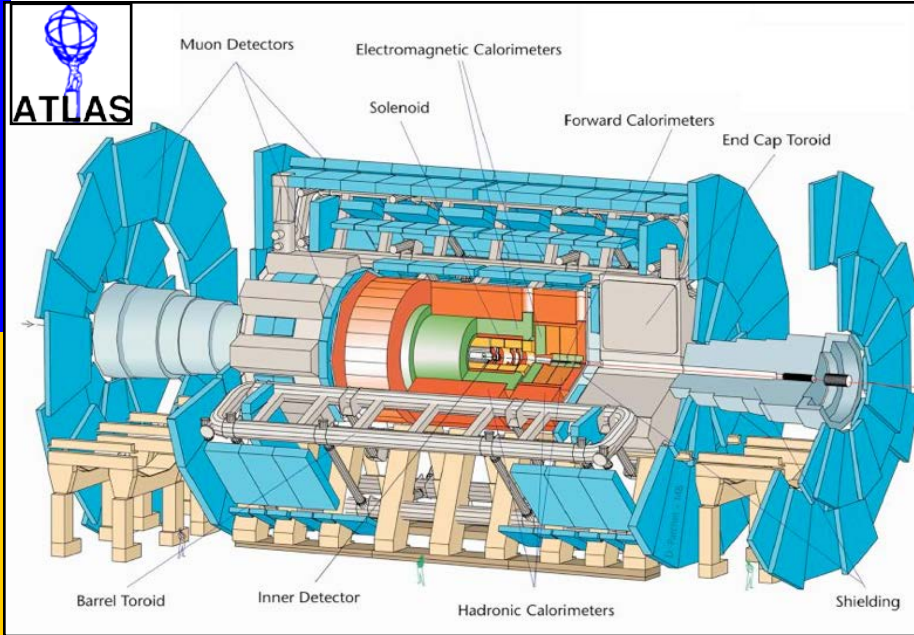
- Relativistic Heavy Ion Collider (RHIC):

- Macroscopic liquid:
  - System size > mean free path
  - System lifetime > relaxation time
- Perfect: shear viscosity/entropy  $\sim 0$

- LHC :

- System is bigger, denser, hotter
- Abundant production of hard probes

# LHC Heavy Ion Program

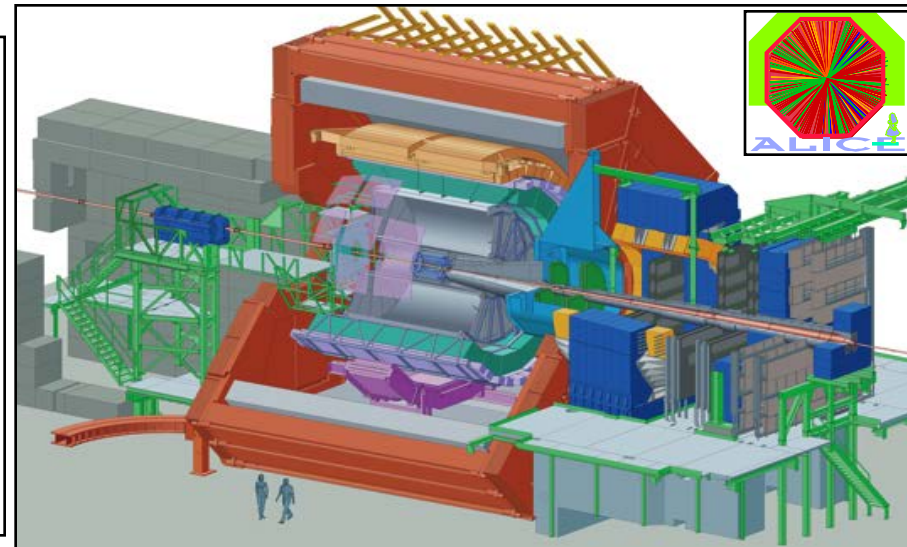
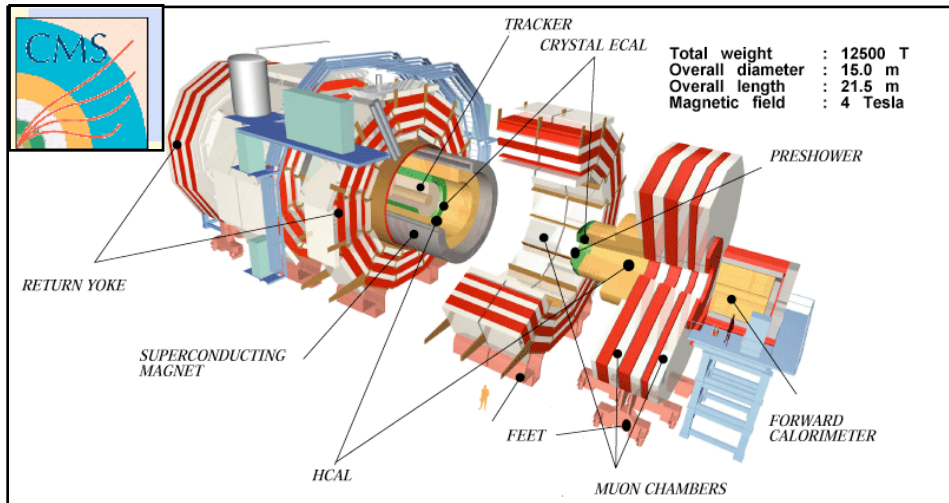


## LHC Heavy Ion Data-taking

Design: Pb + Pb at  $\sqrt{s_{NN}} = 5.5$  TeV  
(1 month per year)

Nov. 2010: Pb + Pb at  $\sqrt{s_{NN}} = 2.76$  TeV

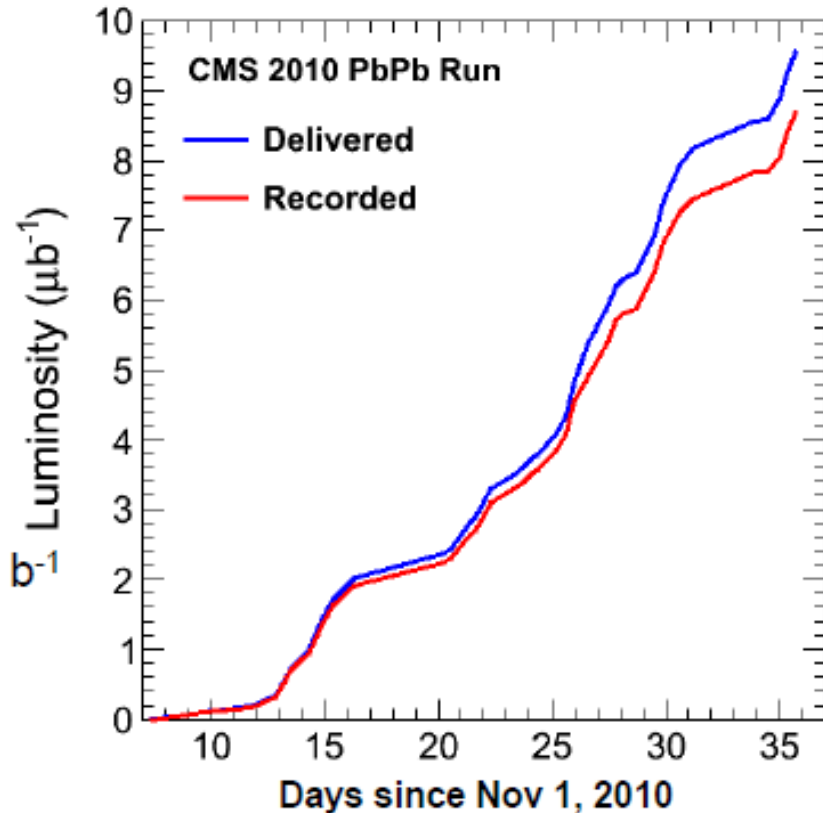
- LHC Collider Detectors
  - ATLAS
  - CMS
  - ALICE





# Pb–Pb Luminosity

*B. Wyslouch, CMS, EPIC2011*



Delivered integrated luminosity  $\sim 9 \mu\text{b}^{-1}$

Luminosity achieved  
 $L = 2\text{--}3 \times 10^{25} \text{ cm}^{-2}\text{s}^{-1}$

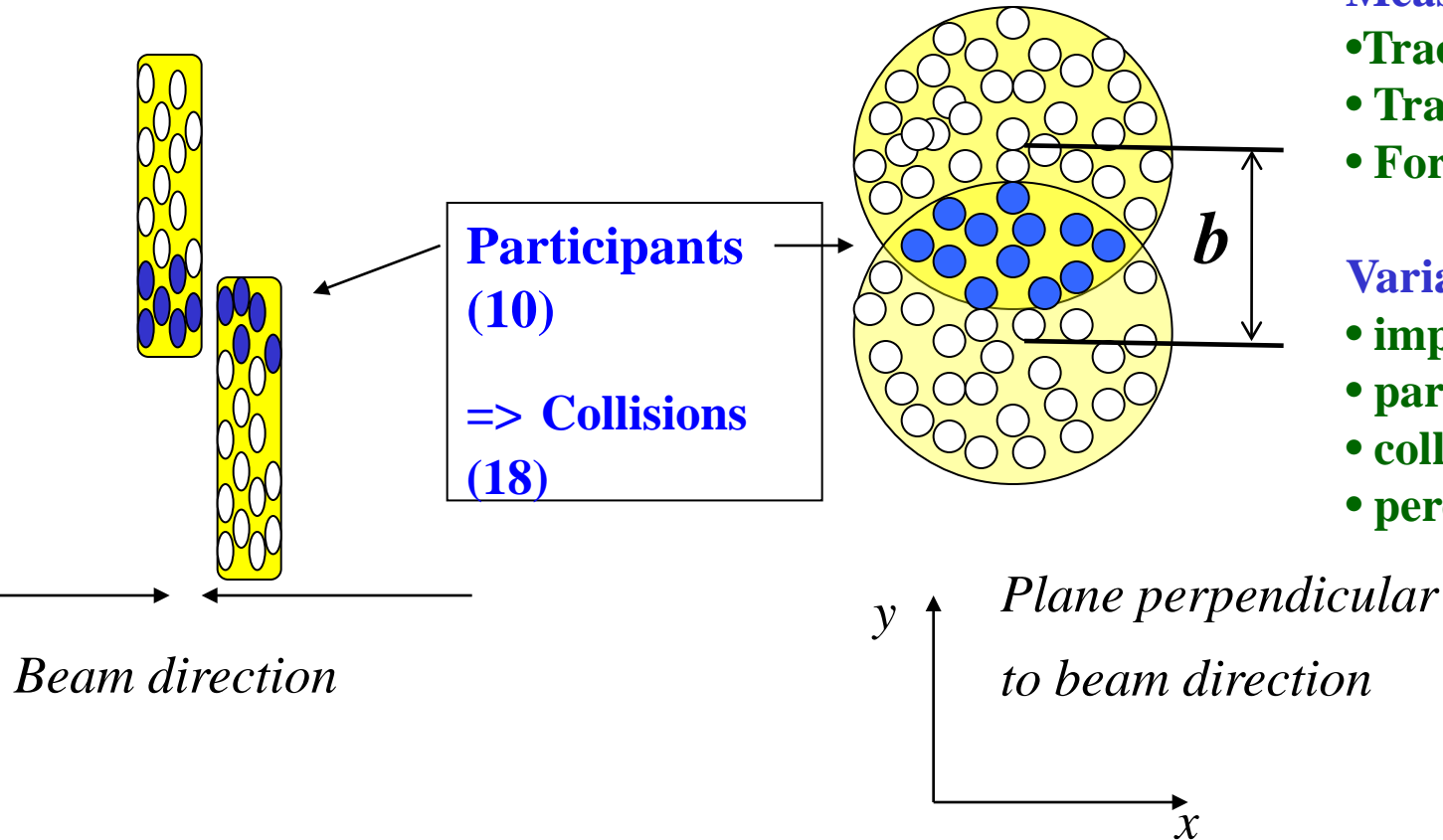
ATLAS very similar to CMS

ALICE recorded  $\sim 50\%$   
due to TPC dead time

# Heavy Ion Collision Centrality

Controls the volume and shape of the system

Multiplicity and energy of produced particles are correlated with geometry of collisions.



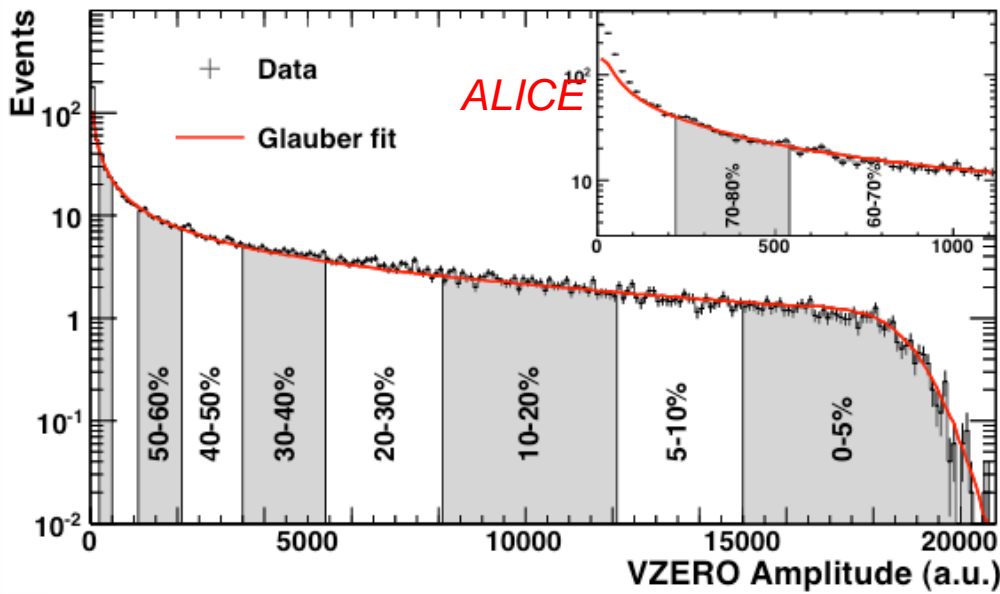
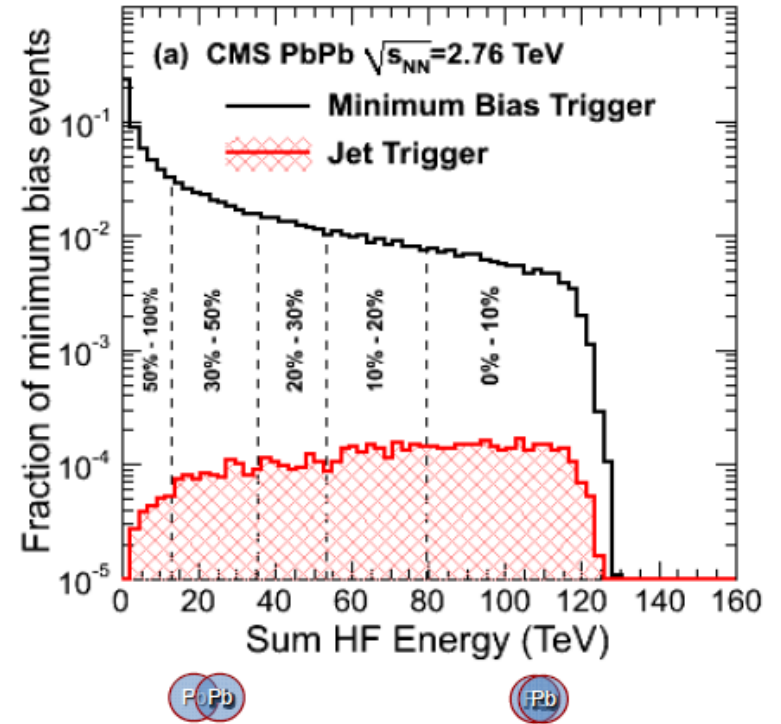
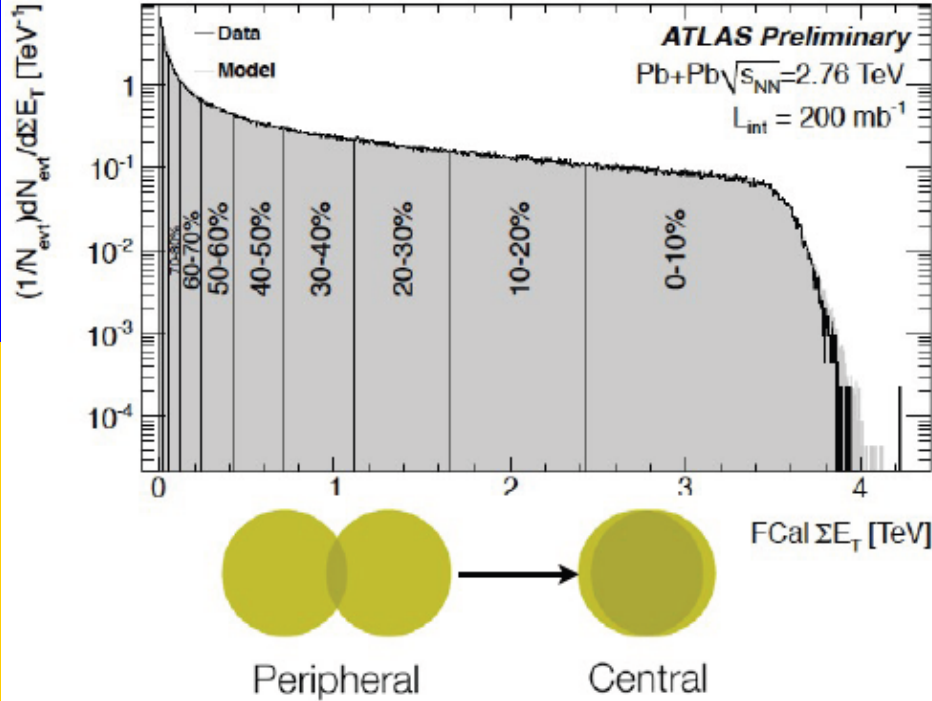
**Measured distribution:**

- Track multiplicity
- Transverse energy
- Forward energy

**Variables:**

- impact parameter
- participants
- collisions
- percentile of x section

# Centrality selection



S.White, ATLAS, EPIC2011

B.Wyslouch, CMS, EPIC2011

M.Nicassio, ALICE, EPIC2011

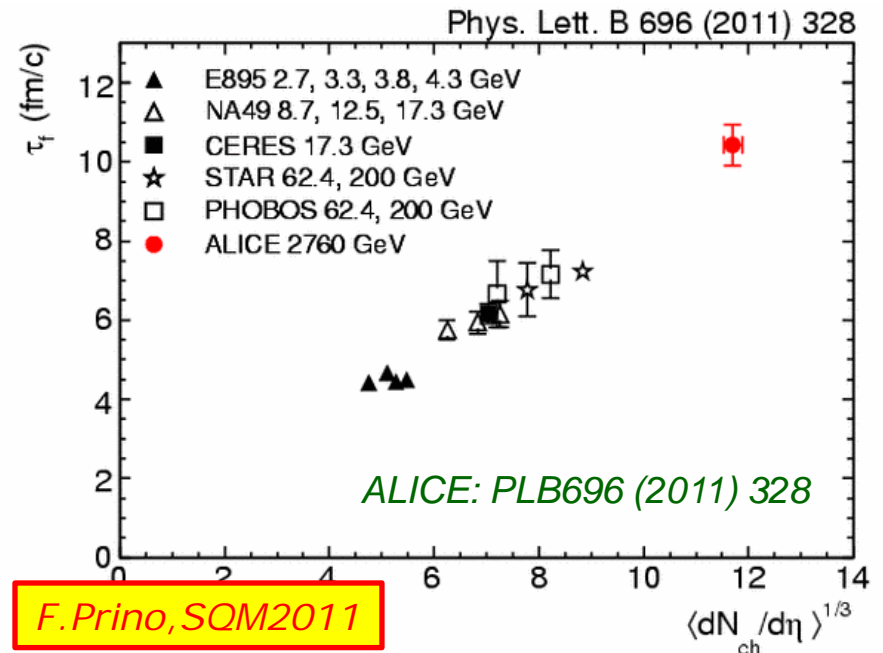
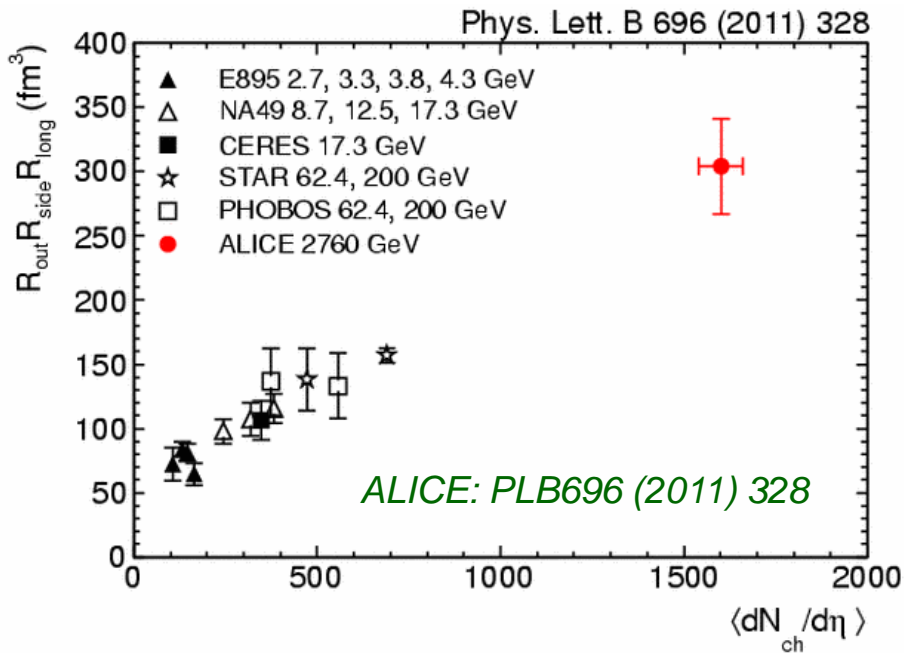
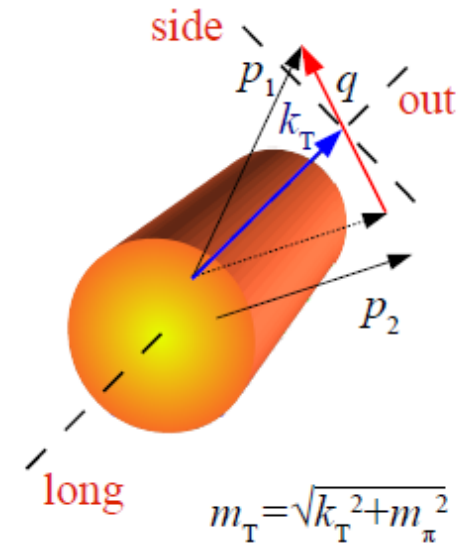


# Soft Probes

- Interferometry of identical particles
- Charged particle multiplicity ,  $E_T$ ,  $\varepsilon$
- Transverse momentum spectra
- Radial flow
- Anisotropic flow

# System size

- Spatial extent of the particle emitting source extracted from interferometry of identical bosons
  - Two-particle momentum correlations in 3 orthogonal directions -> HBT radii ( $R_{\text{long}}$ ,  $R_{\text{side}}$ ,  $R_{\text{out}}$ )
  - Size: twice w.r.t. RHIC
  - Lifetime: 40% higher w.r.t. RHIC



# Multiplicity, $E_T$ and $\varepsilon$

## Particle Production and Energy density $\varepsilon$ :

Produced Particles:  $dN_{ch}/d\eta \approx 1600 \pm 76$  (syst)

$\approx 30,000$  particles in total,  $\approx 400$  times pp !

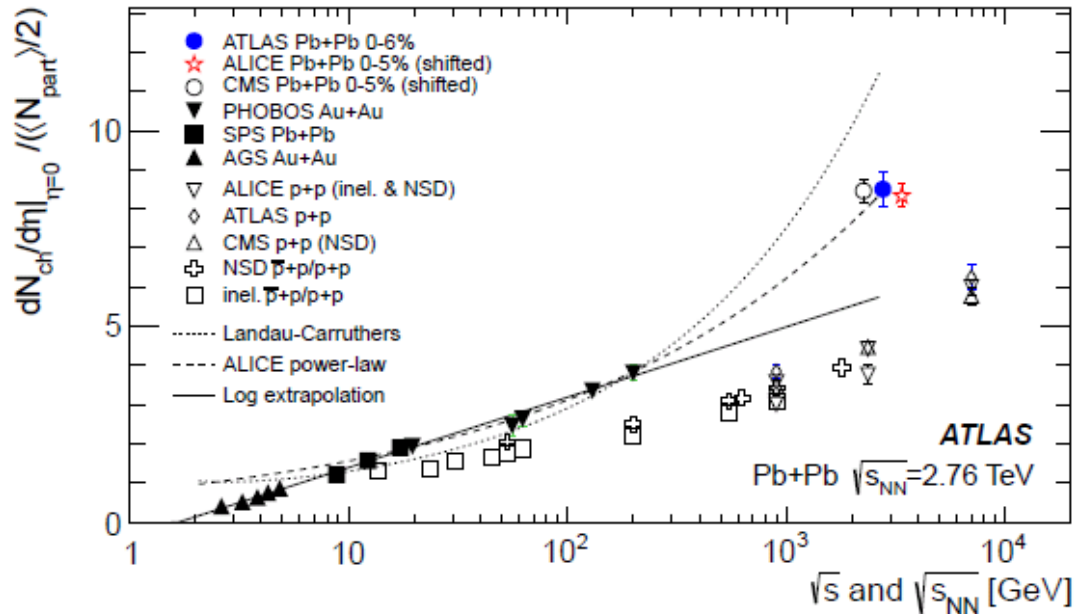
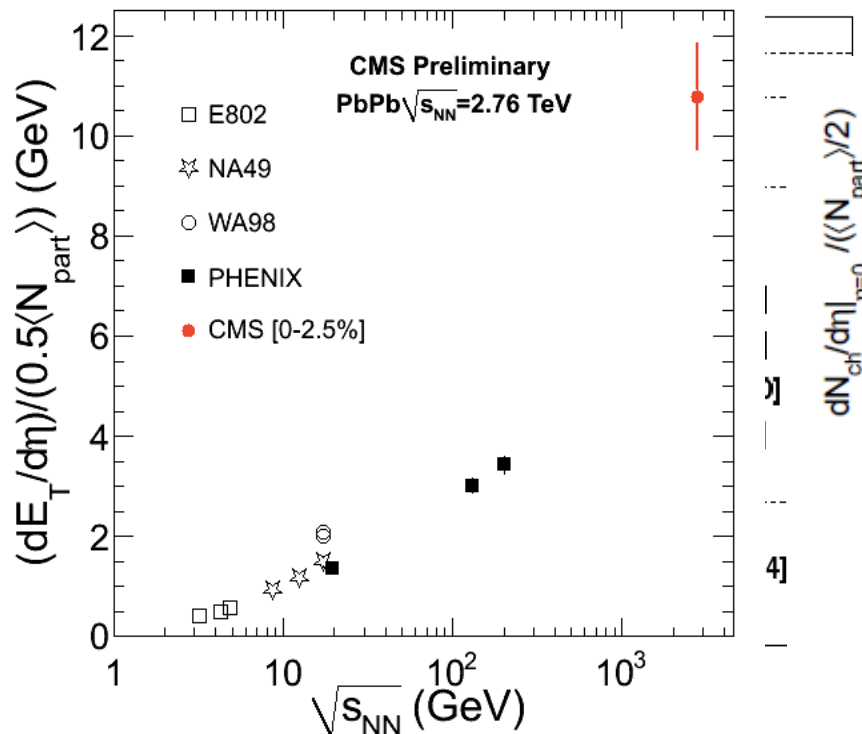
somewhat on high side of expectations (tuned to RHIC)

growth with energy faster in AA

Energy density  $\varepsilon > 3 \times$  RHIC (fixed  $\tau_0$ )

Temperature + 30%

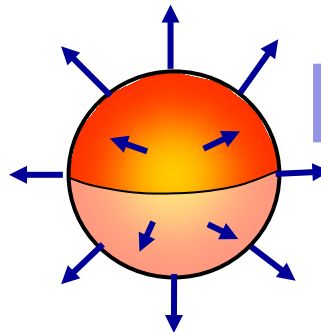
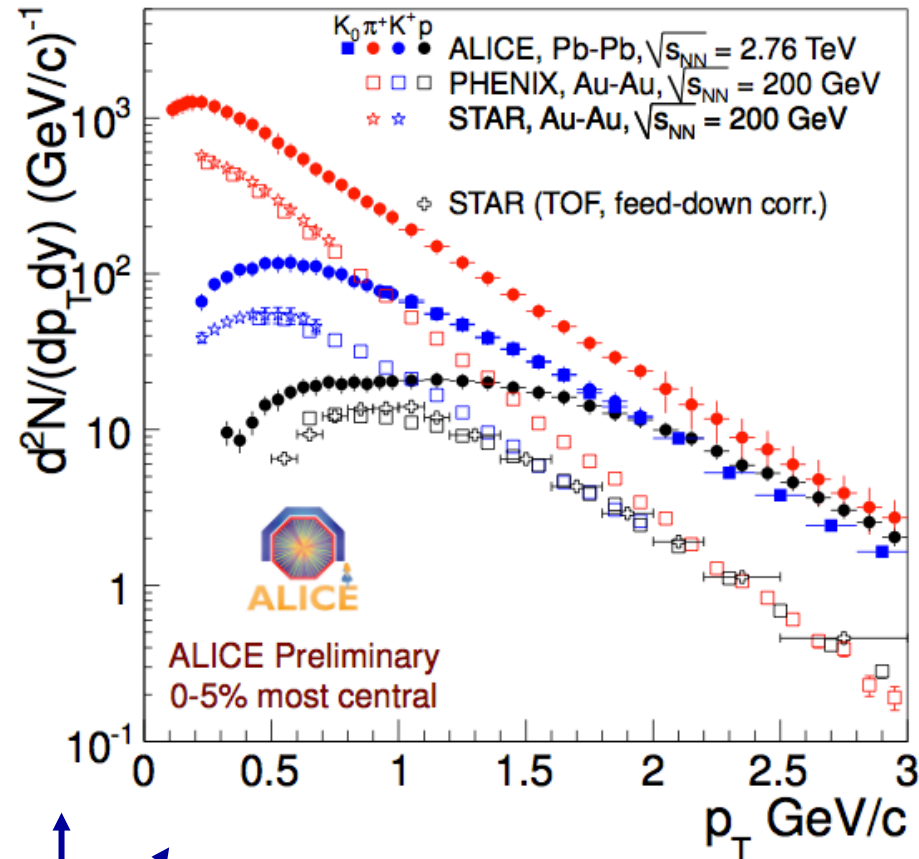
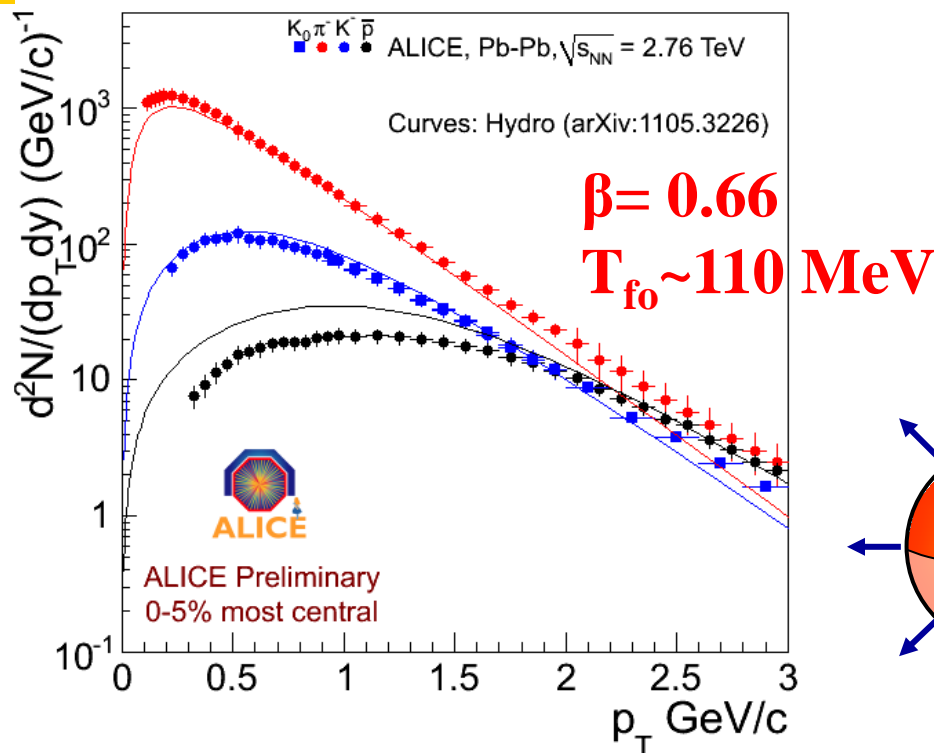
$$\varepsilon(\tau) = \frac{E}{V} = \frac{1}{\tau_0 A} \frac{dE_T}{dy}$$



# Charged particle spectra

## Radial Flow

- K,  $\pi$ , p spectra **0-5% central collisions**
- Very clear flattening and higher tails at  $\sqrt{s_{NN}}=2.76$  TeV
- Quantify with blastwave parameter studies: **radial flow**  $\beta=v_0/c$  and freezeout temperature  $T_{fo}$



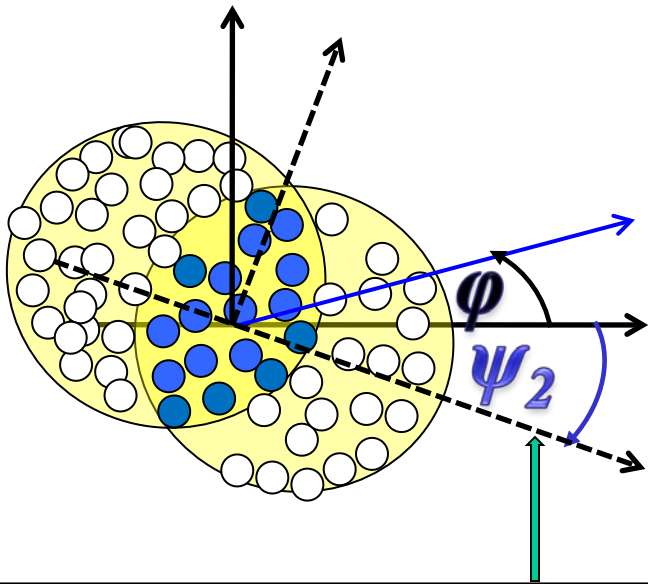
L. Barnby, ALICE, EPIC2011

# Anisotropic Flow

- Fourier expansion in azimuthal distribution:

$$\frac{dN}{p_T dp_T dy d\varphi} = \frac{1}{2\pi} \frac{dN}{p_T dp_T dy} (1 + 2v_1 \cos(\varphi - \psi_1) + 2v_2 \cos(2(\varphi - \psi_2)) + \dots)$$

- $\varphi$  – azimuthal angle



In non-central collisions participant area is not azimuthally symmetric: system evolution transfer this anisotropy from coordinate space to momentum space

$v_1$  - direct flow

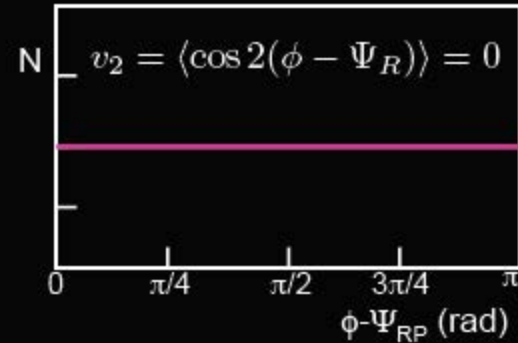
$v_2$  - elliptic flow, dominant for system symmetric wrt Collision Plane<sub>12</sub>

Collision Plane :  
- Defined by Beam and Impact Parameter

# Elliptic flow - $v_2$

## 1) superposition of independent n+n:

momenta pointed at random  
relative to reaction plane

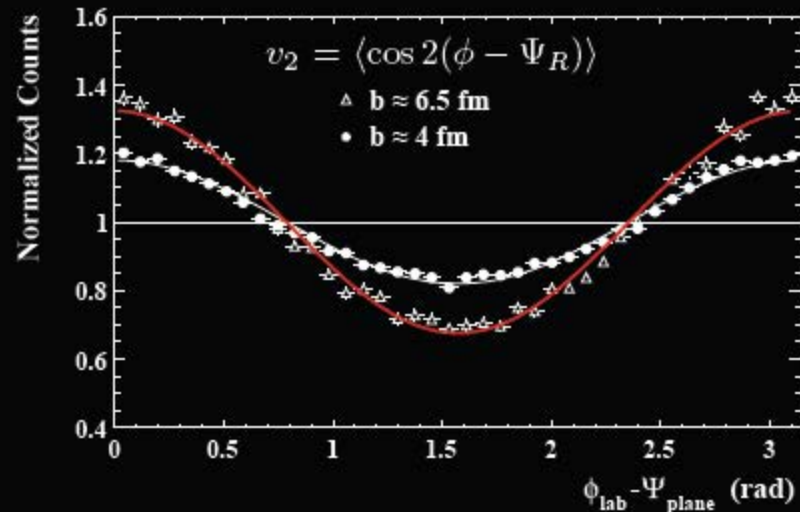


## 2) evolution as a bulk system

pressure gradients (larger in-plane)  
push bulk "out"  $\rightarrow$  "flow"



more, faster particles  
seen in-plane



nonzero  $v_2$  indicates a departure from independent n+n!



# Physics of elliptic flow

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Elliptic flow depends on:

- Initial conditions
- Fluid properties
  - Equation of state
  - Shear viscosity

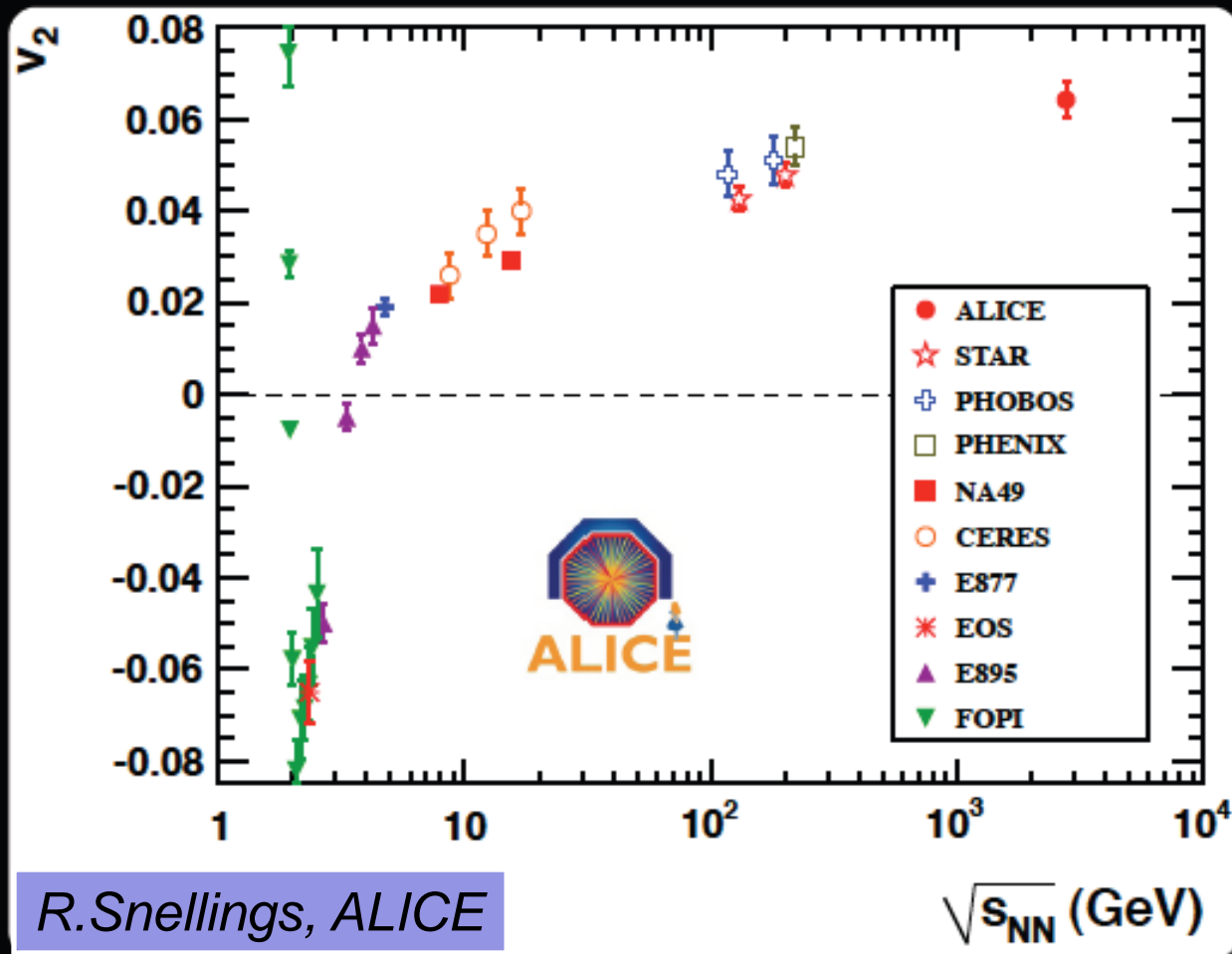
Shear viscosity:  $\eta = n \langle p \rangle \lambda = \langle p \rangle / \sigma$

*Small viscosity  $\eta \Rightarrow$  large cross section  $\sigma$   
 $\Rightarrow$  strongly interacting fluid*

# The Perfect Liquid



K. Aamodt et al. (ALICE Collaboration)  
PRL 105, 252302 (2010)

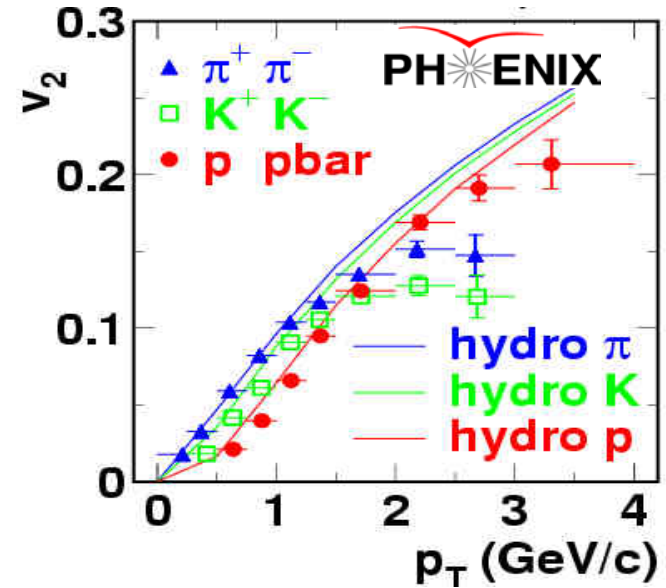
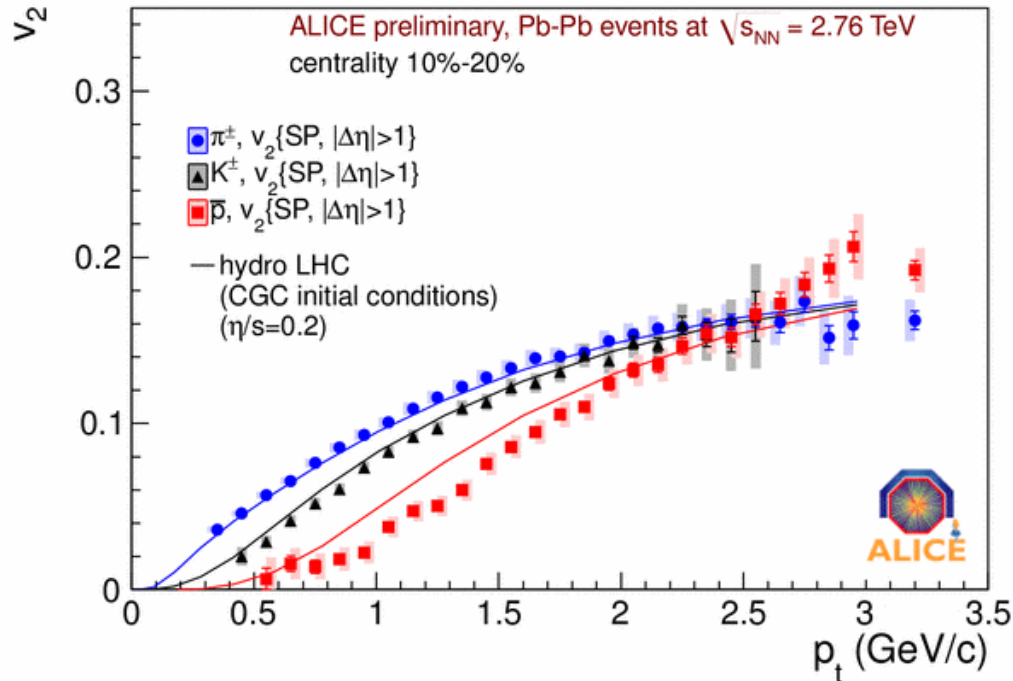


R. Snellings, ALICE

The system produced at the LHC behaves as a very low viscosity fluid (a perfect fluid)

# Hydrodynamics and $v_2$

- comparison of identified particles  $v_2(p_T)$  with hydro prediction – mass splitting described

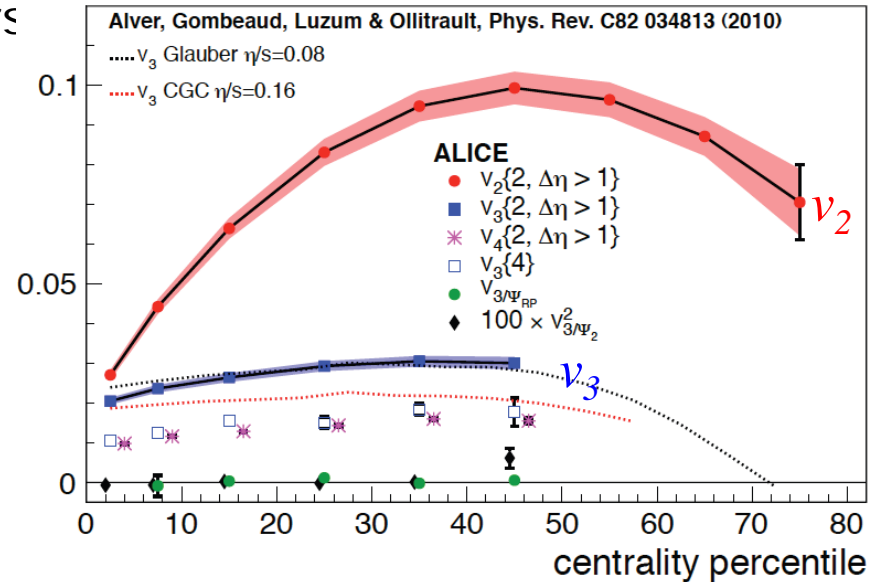
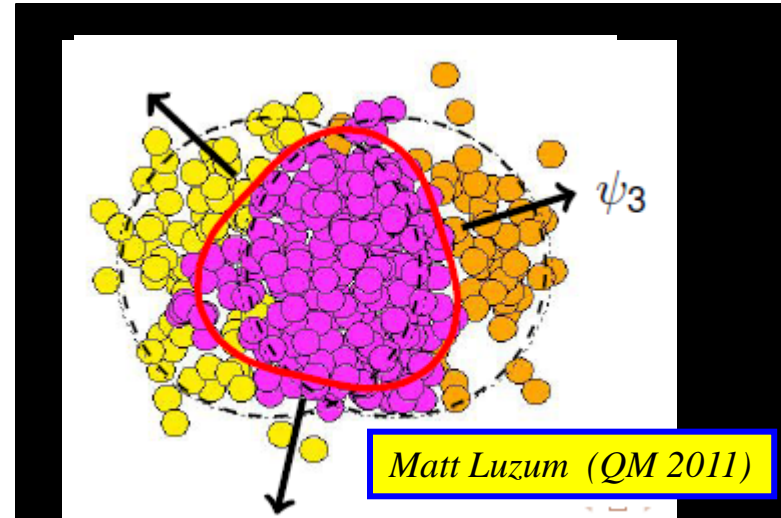


ALI-PREL-2448

- (calculation by C Shen et al.: arXiv:1105.3226 [nucl-th])
- Protons are to be understood**

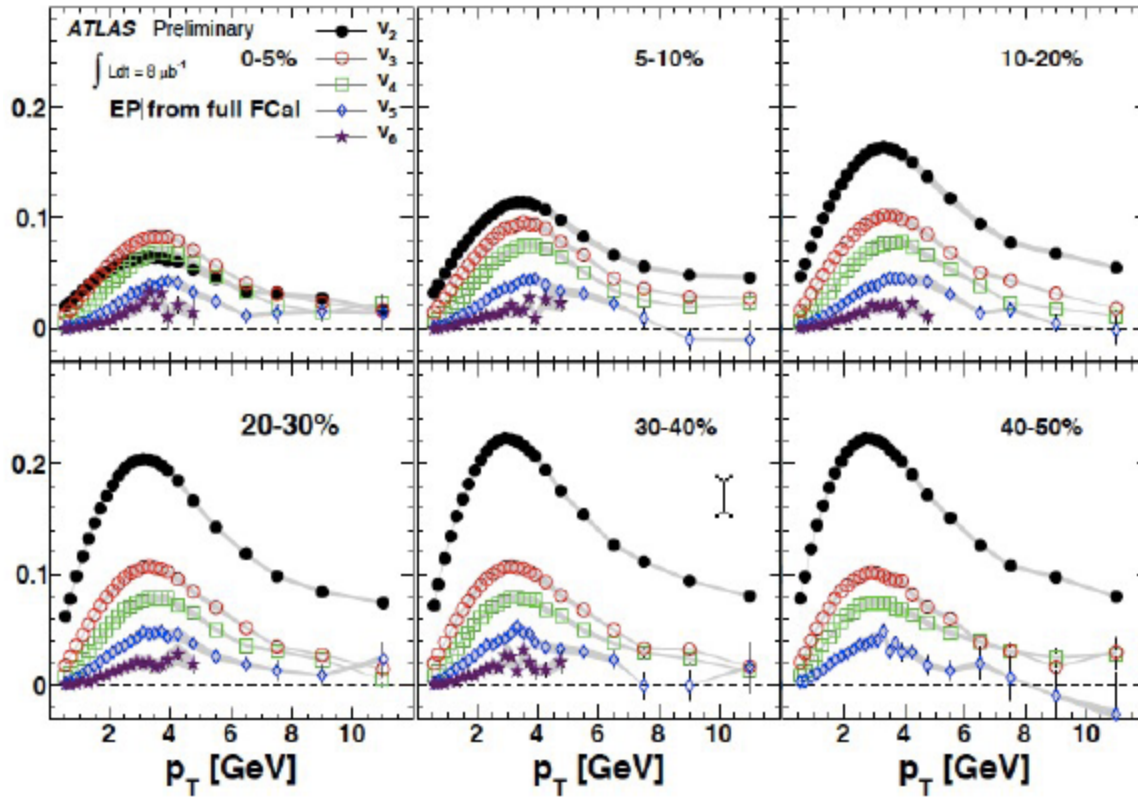
# Fluctuations $\rightarrow v_3$

- “ideal” shape of participants’ overlap is  $\sim$  elliptic
  - in particular: no odd harmonics expected
  - participants’ plane coincides with event plane
- but fluctuations in initial conditions:
  - participants plane  $\neq$  event plane
  - $\rightarrow v_3$  (“triangular”) harmonic appears
- [B Alver & G Roland, PRC81 (2010) 054905]
- and indeed,  $v_3 \neq 0$  !
- $v_3$  has weaker centrality dependence than  $v_2$



ALICE: PRL 107 (2011) 032301

# Higher harmonics



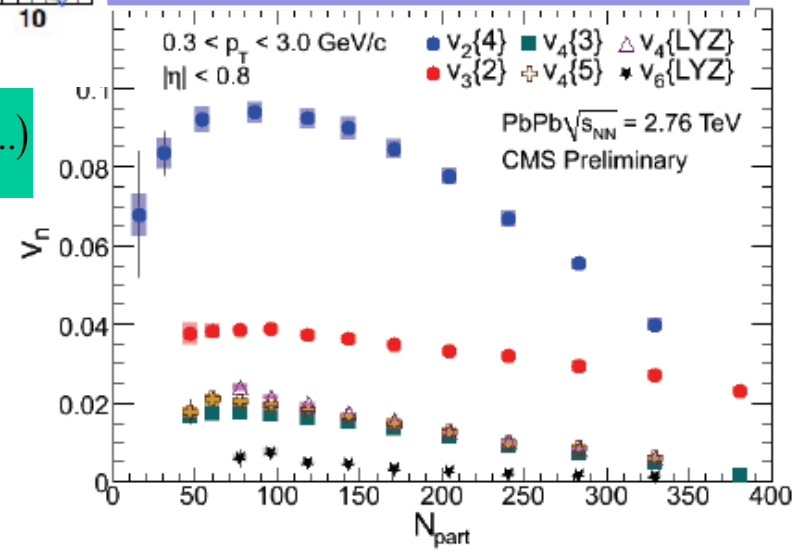
*S. White, ATLAS, EPIC2011*

- $v_{n+1} < v_n$
- $v_{n+1}$  less centrality dependent than  $v_n$

$$\frac{dN}{p_T dp_T dy d\phi} = \frac{1}{2\pi} \frac{dN}{p_T dp_T dy} (1 + 2v_2 \cos(\phi - \psi_2) + 2v_3 \cos(2(\phi - \psi_3)) + \dots)$$

$v_n$  – information on viscosity,  $n > 2$

*M. Issah, CMS, EPIC2011*



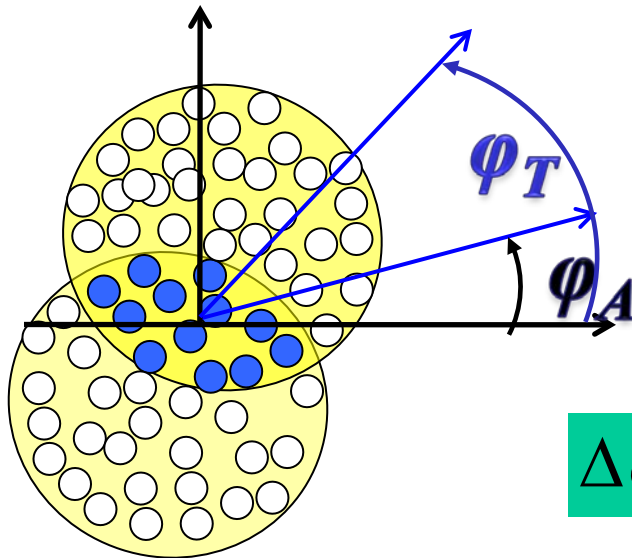
# 2 Particle Correlations and Flow

- Fourier expansion in azimuthal distribution:

$$\frac{dN}{d\Delta\varphi} = (1 + 2v_1 \cos(\Delta\varphi) + 2v_2 \cos(2\Delta\varphi) + \dots)$$

- If flow dominates than:

$$v_i = v_i^A * v_i^T$$



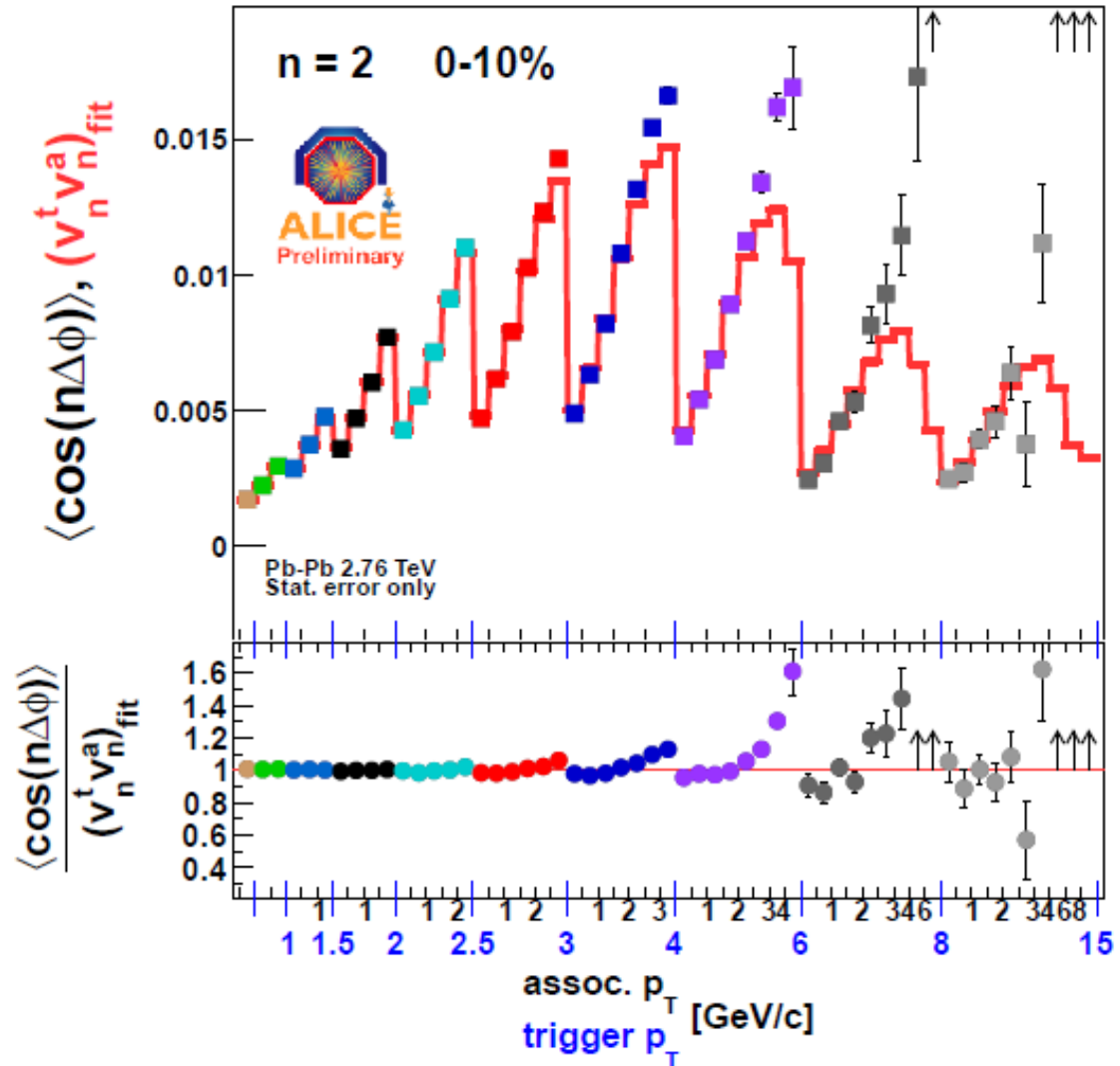
$$\Delta\varphi = \varphi_T - \varphi_A$$

Flow Fourier coefficients



# Flow vs Non-Flow Correlations

- Compare single calculated values with global fit
- To some extent, a good fit suggests flow-type correlations, while a poor fit implies non-flow effects
- $v_2$  to  $v_5$  factorize until  $p_T \sim 3-4$  GeV/c, then jet-like correlations dominate
- $v_1$  factorization problematic (influence of away-side jet)

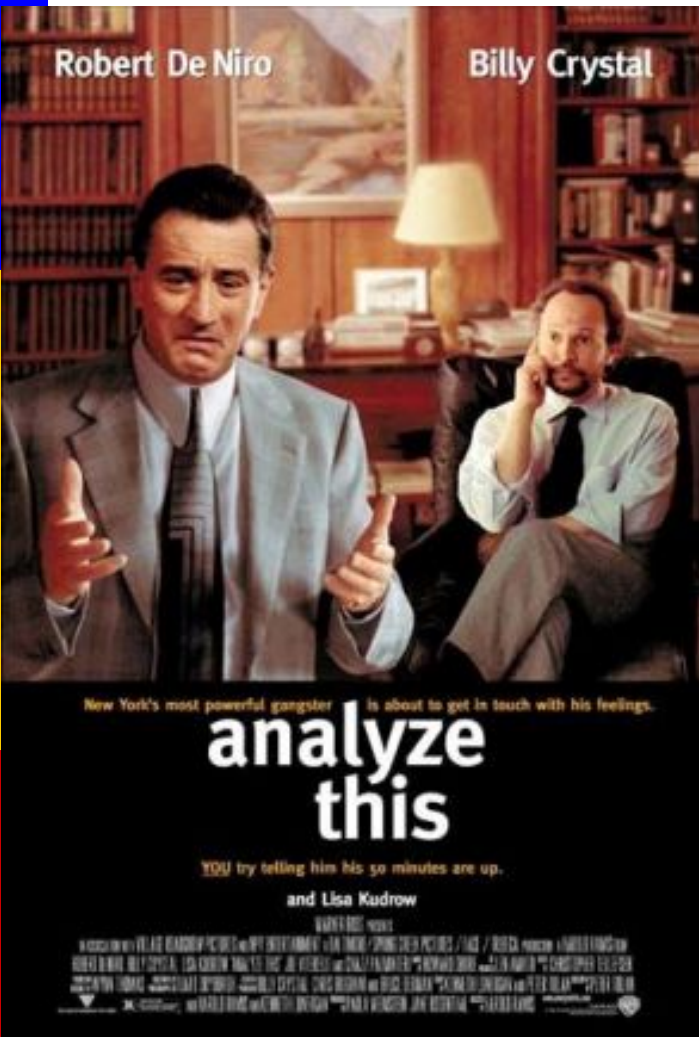


# Anisotropic Flow Summary

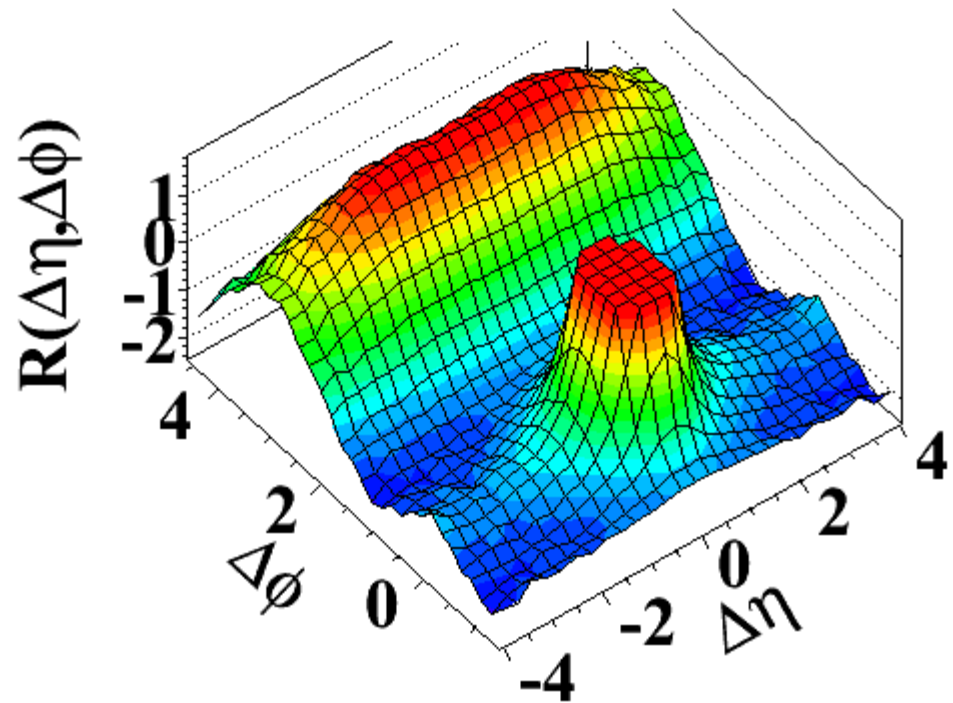
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- Centrality and  $p_t$  dependences of various  $v_n$  constraint
  - initial conditions (CGC vs Glauber)
  - viscosity –  $\eta/s$
- There is no hydro calculation (yet) describing simultaneously data on  $v_2$  and  $v_3, \dots$
- 2 particle correlations consistent with flow for  $p_T < 3-4$  GeV/c

# Speaking of which...



(d)  $N > 110$ ,  $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$



Full Fourier decomposition of the CMS pp ridge?

# The nuclear modification factor

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- quantify departure from binary scaling in AA
  - ratio of yield in AA versus reference collisions
- e.g.: reference is pp  $\rightarrow R_{AA}$

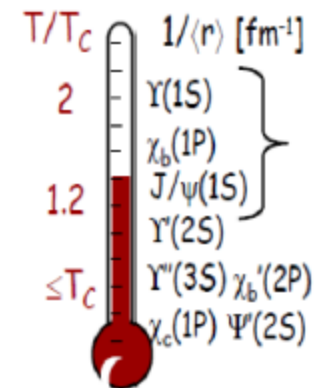
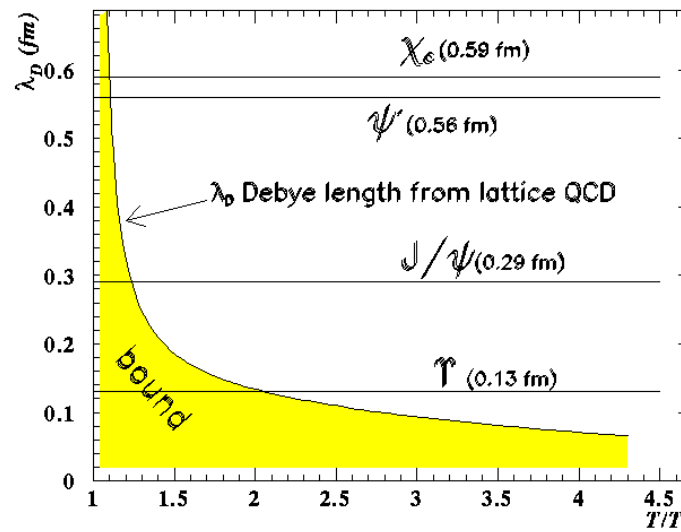
$$R_{AA} = \frac{\text{Yield}_{AA}}{\text{Yield}_{pp}} \cdot \frac{1}{\langle Nbin \rangle_{AA}}$$

- ...or peripheral AA  $\rightarrow R_{CP}$  ("central to peripheral")

$$R_{cp} = \frac{\text{Yield}_{AA, \text{central}}}{\text{Yield}_{AA, \text{periph}}} \cdot \frac{\langle Nbin \rangle_{AA, \text{periph}}}{\langle Nbin \rangle_{AA, \text{central}}}$$

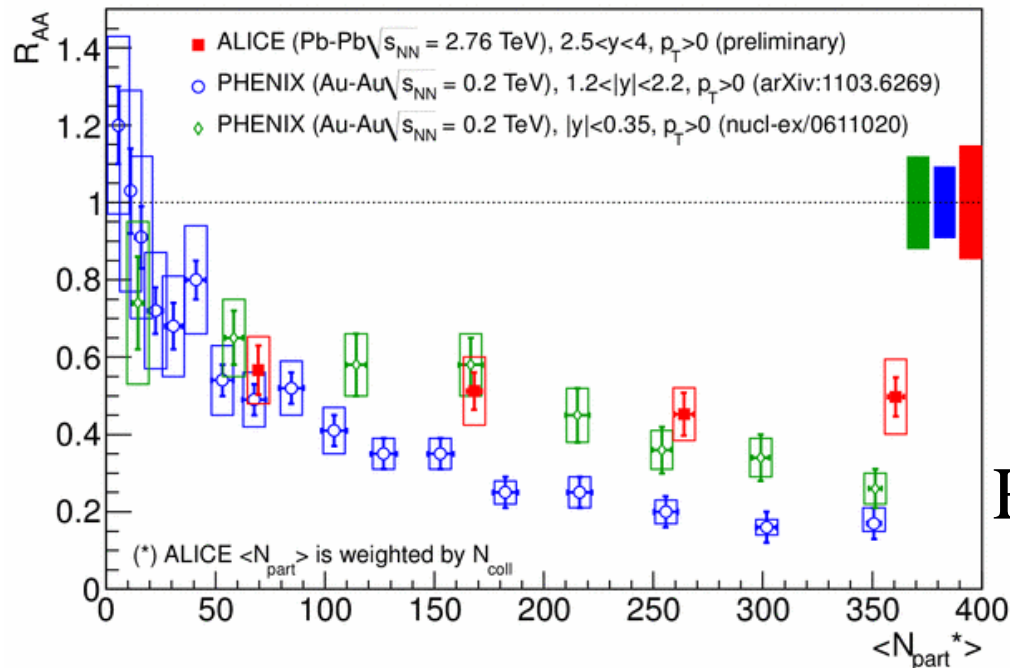
# Quarkonia suppression

- In the plasma phase the interaction potential is expected to be screened beyond the Debye length  $\lambda_D$  (analogous to e.m. Debye screening):
- Charmonium (cc) and bottomonium (bb) states with  $r > \lambda_D$  will not bind; their production will be **suppressed**
- Recombination of cc and bb **regenerates** quarkonia



# J/ψ @ LHC: forward y, low p<sub>T</sub>

- LHC: 2.5 < y < 4, p<sub>T</sub> > 0 (ALICE)
- Less suppression than RHIC:  
1.2 < y < 2.2, p<sub>T</sub> > 0 (PHENIX)
- As suppressed as RHIC: |y| < 0.35, p<sub>T</sub> > 0 (PHENIX)



$$R_{AA} = \frac{\text{Yield}_{AA}}{\text{Yield}_{pp}} \cdot \frac{1}{\langle N_{bin} \rangle_{AA}}$$

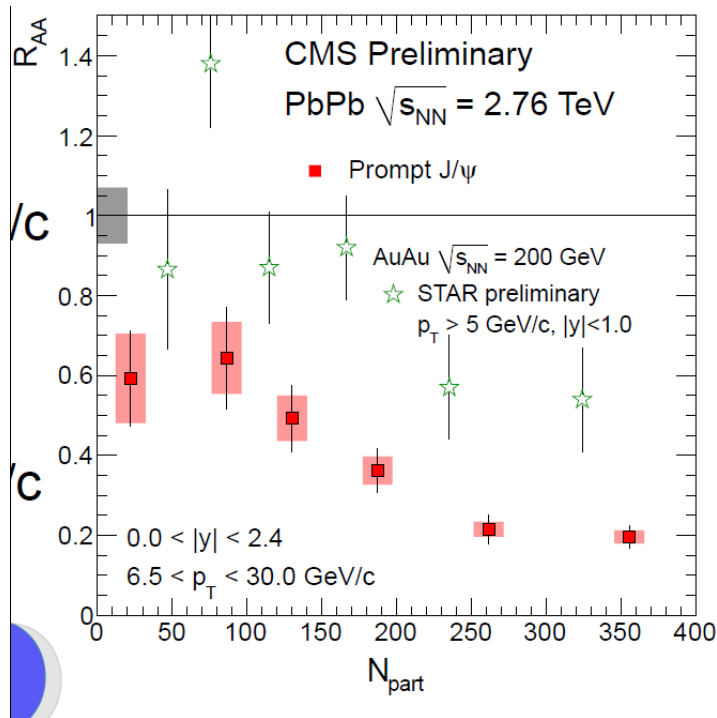
Recombination ?



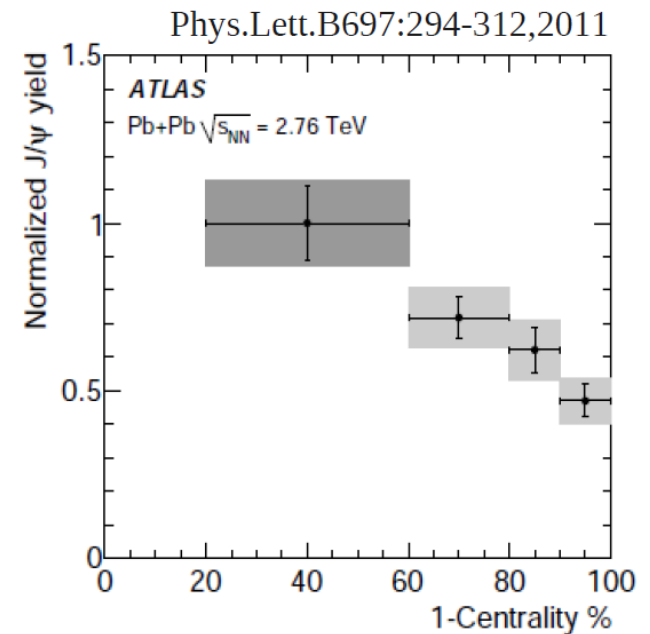
# J/ψ @ LHC: central y, high p<sub>T</sub>

- LHC:  $|y| < 2.4$ ,  $p_T > 6.5$  GeV/c (CMS) prompt J/ψ

$$R_{AA} = \frac{\text{Yield}_{AA}}{\text{Yield}_{pp}} \cdot \frac{1}{\langle N_{bin} \rangle_{AA}}$$



CMS: PAS HIN-10-006

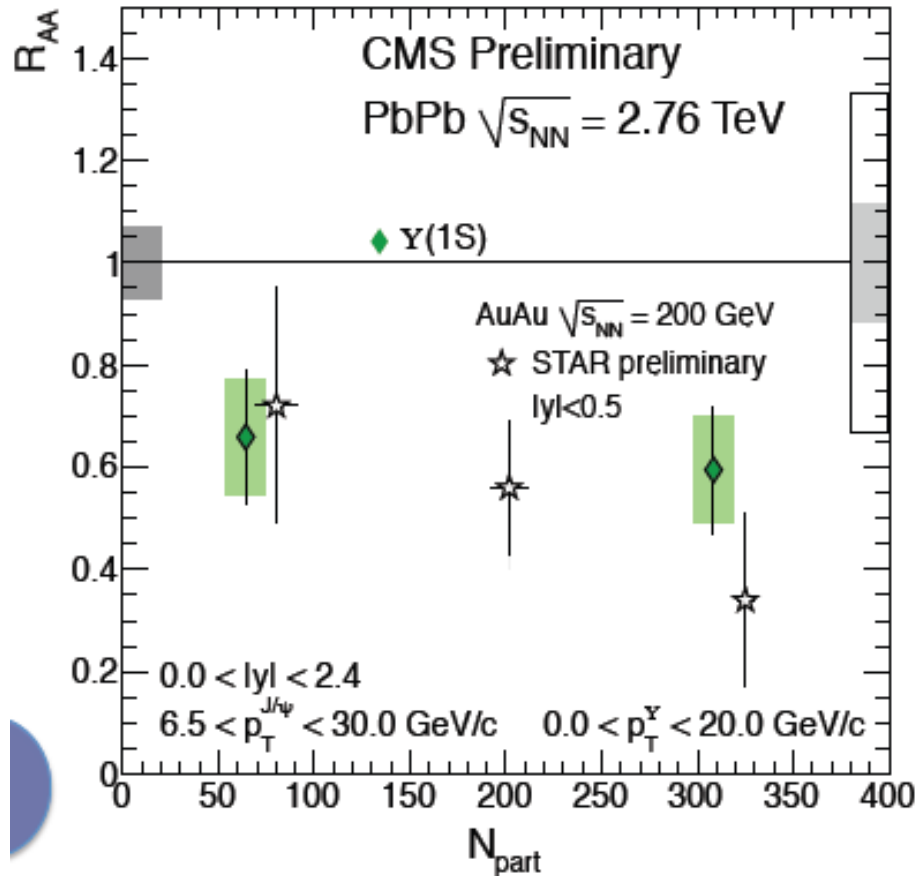


ATLAS: PLB 697 (2011) 294

- more suppressed than RHIC:
- $|y| < 1$ ,  $p_T > 5$  GeV/c (STAR)  
inclusive J/ψ

# $\Upsilon(1S)$ suppression

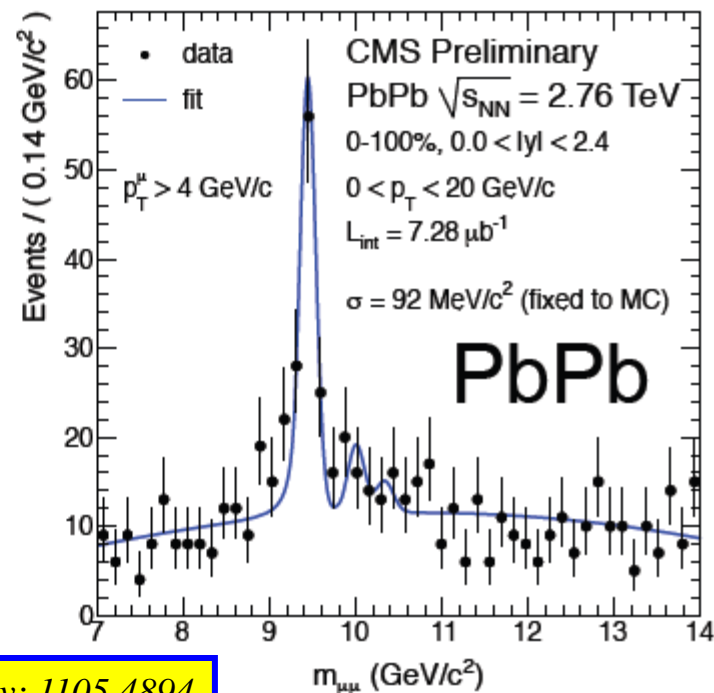
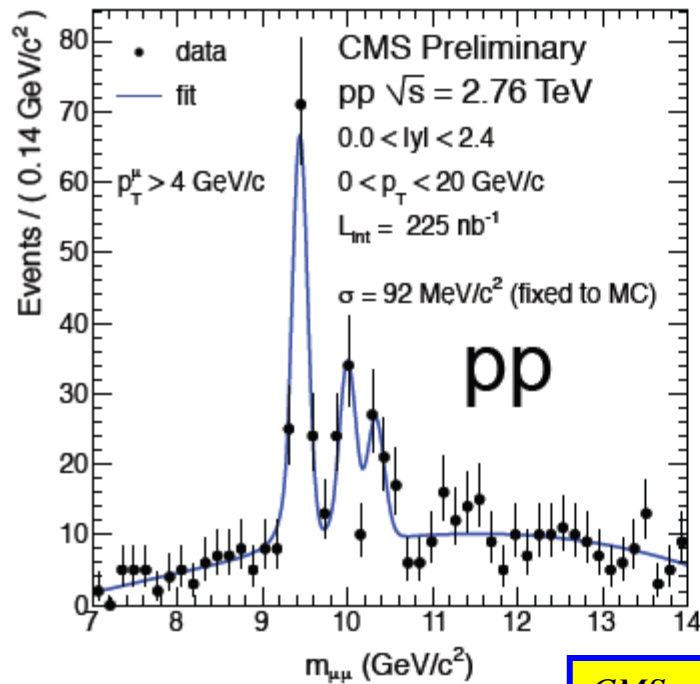
CMS: PAS HIN-10-006



- CMS  $\Upsilon(1S)$   $R_{AA}(0-100) = 0.62 \pm 0.11 \pm 0.10$
- STAR  $\Upsilon(1+2+3S)$   $R_{AA}(0-60) = 0.56 \pm 0.11^{+0.02}_{-0.10}$

# $\Upsilon(2S+3S)$ suppression

- additional suppression for  $\Upsilon(2S+3S)$  w.r.t.  $\Upsilon(1S)$  ?



CMS: arXiv: 1105.4894

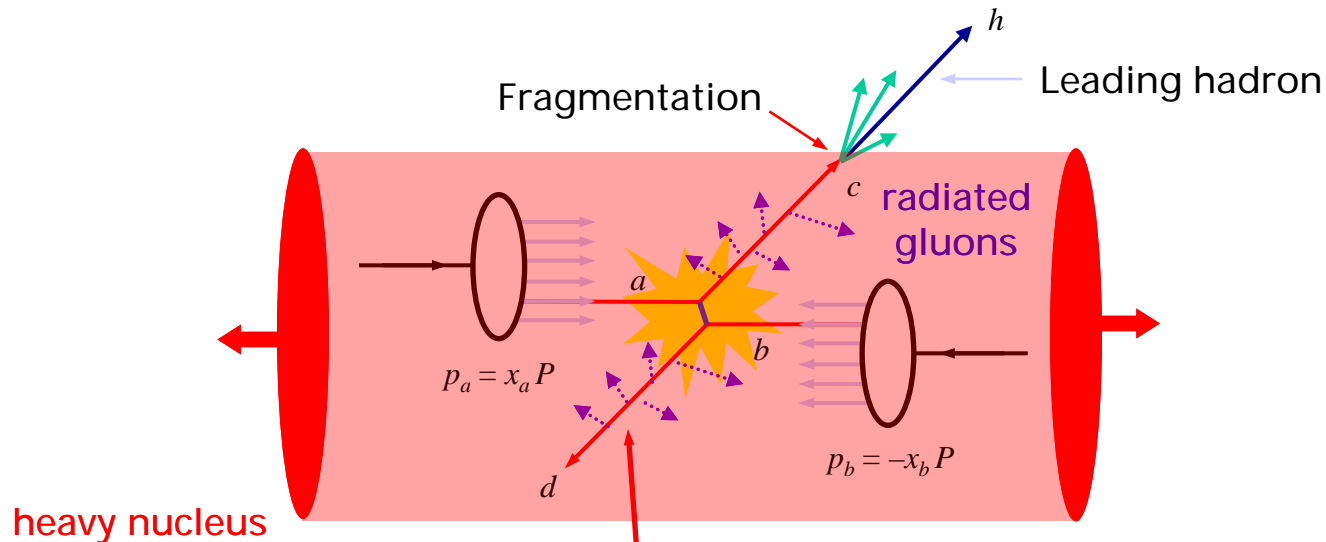
$$\frac{\Upsilon(2S + 3S)/\Upsilon(1S)|_{\text{Pb-Pb}}}{\Upsilon(2S + 3S)/\Upsilon(1S)|_{pp}} = 0.31^{+0.19}_{-0.15}(\text{stat}) \pm 0.03(\text{syst}),$$

# Quarkonia Summary

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- Y and  $J/\psi$  suppressed by same amount ?
- Suppression depends on  $y$  and  $p_t$
- the future runs should allow us to establish quantitatively the complete quarkonium suppression(/recombination?) pattern
  - high statistic measurements
  - open flavour baseline / contamination
  - pA baseline

# Jets in medium



**Key prediction:** jets are quenched

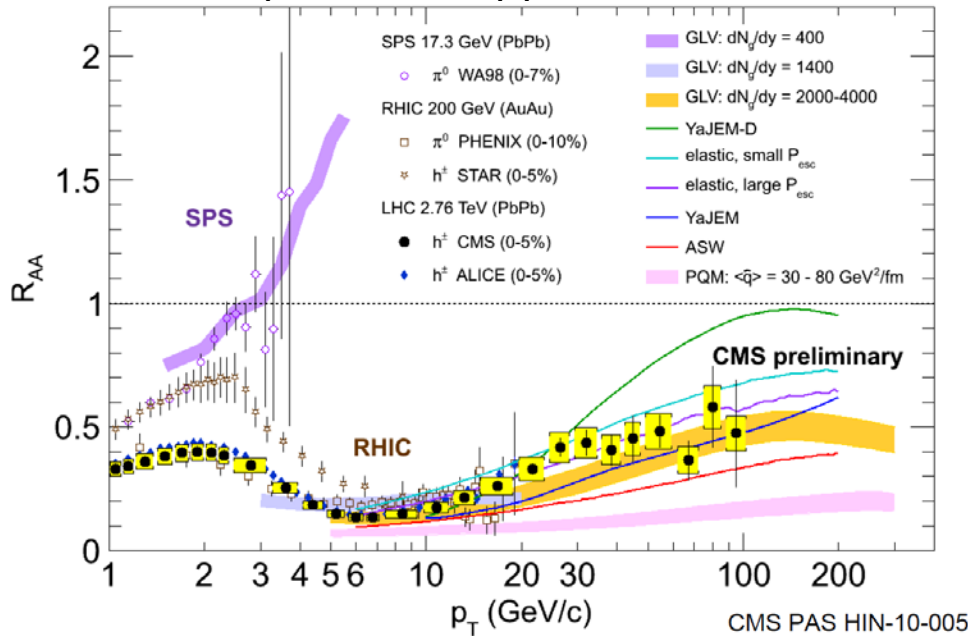
- collisional energy loss (Bjorken)
- radiative energy loss (Wang and Gyulassy)

$$R_{AA} = \frac{\text{Yield}_{AA}}{\text{Yield}_{pp}} \cdot \frac{1}{\langle N_{bin} \rangle_{AA}}$$

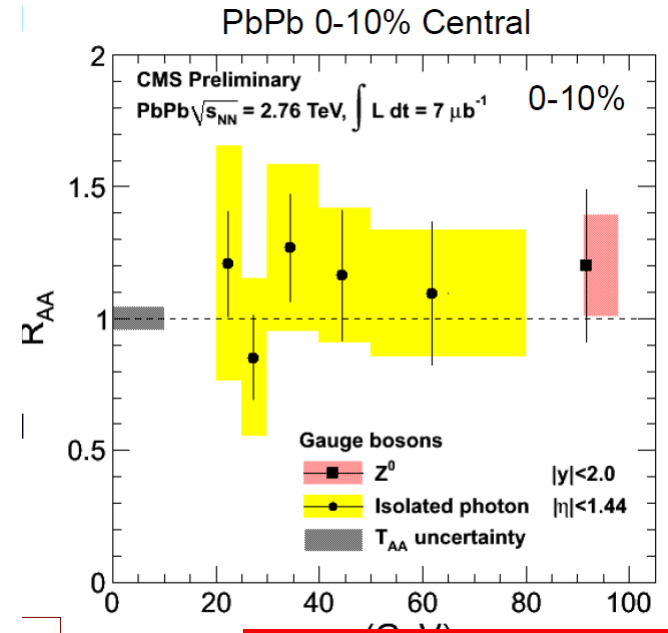
# $R_{AA}$ at LHC

Non colour probes

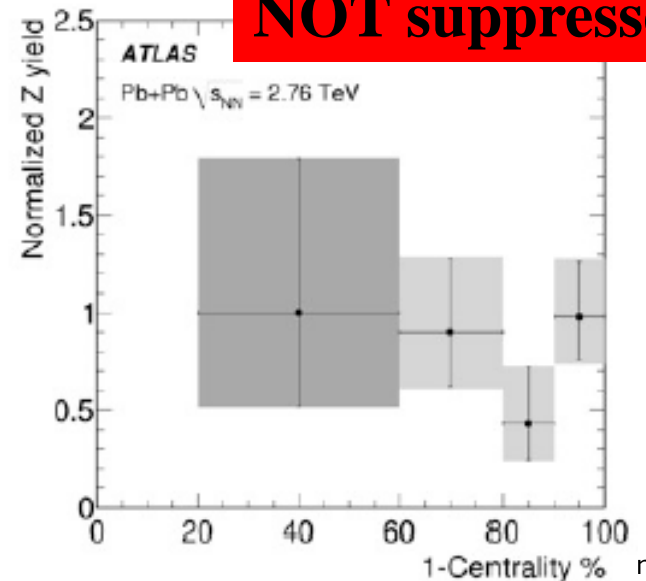
- $R_{AA}(p_T)$  for charged particles produced in 0-5% centrality range
  - minimum ( $\sim 0.14$ ) for  $p_T \sim 6-7$  GeV/c
  - then slow increase at high  $p_T$
  - still significant suppression



- essential quantitative constraint for parton energy loss models!



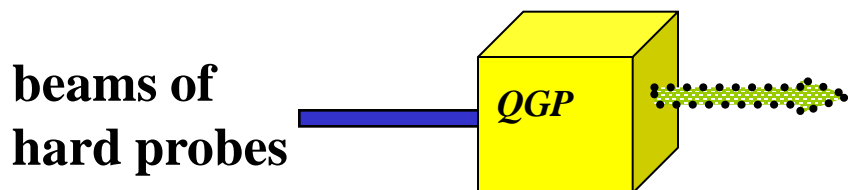
**NOT suppressed**





# Jet quenching

- Jet quenching = Energy loss of fast parton in matter
  - **jet(parton)  $E \rightarrow \text{jet } E' (=E-\Delta E) + \text{soft gluons } (\Delta E)$**
  - **modified jet fragmentation function  $f(z)$**  (= number & energy distribution of hadrons) via matter **induced gluon radiation**/scattering
- QCD energy loss  $\Delta E = f(m) \times c_q \times q \times L^2 \times f(E)$  depends on:
  - **$q$**  : 'transport coefficient ' = property of medium (QGP >> nuclear matter)
  - **$L$** : size of medium ( $\sim L^2$ )
  - **$c_q$** : parton type (gluon > quark)
  - **$f(m)$**  : quark mass (light  $q$  > heavy  $Q$ )
  - **$f(E)$**  : jet energy ( $\Delta E = \text{constant}$  or  $\sim \ln(E)$ )

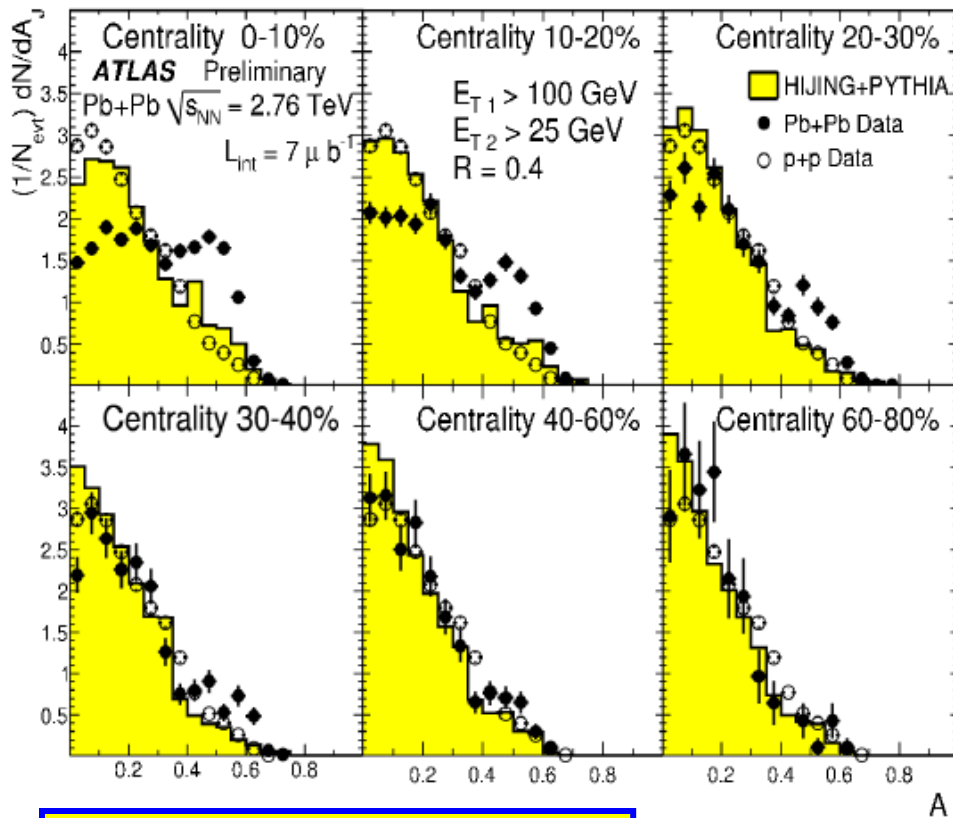


- 1) How much energy is lost ?  
measure 'hard' fragments
- 2) Where (and how) is it lost ?
- 3) Shows expected scaling ?

# How much is lost ?

- imbalance quantified by the di-jet asymmetry variable  $A_J$ :

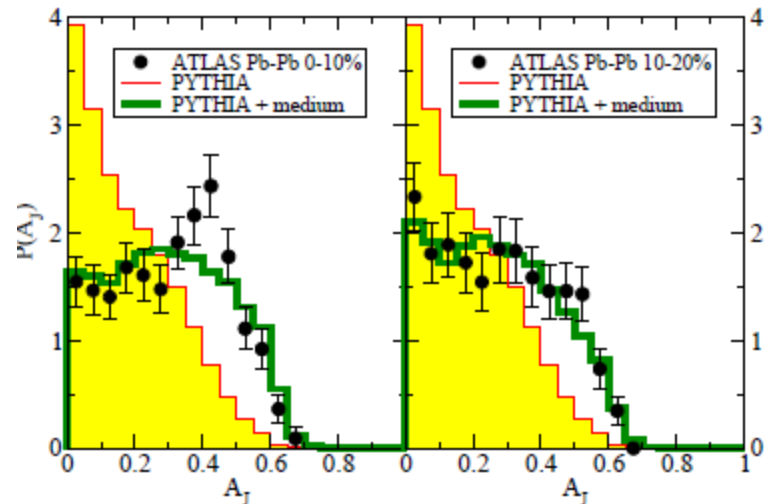
$$R = 0.4 \quad |\eta| < 2.8$$



ATLAS: PRL105 (2010) 252303

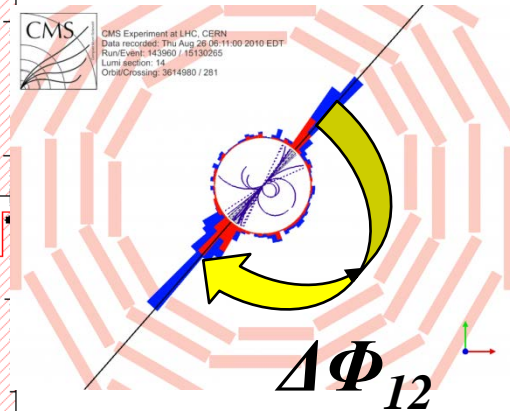
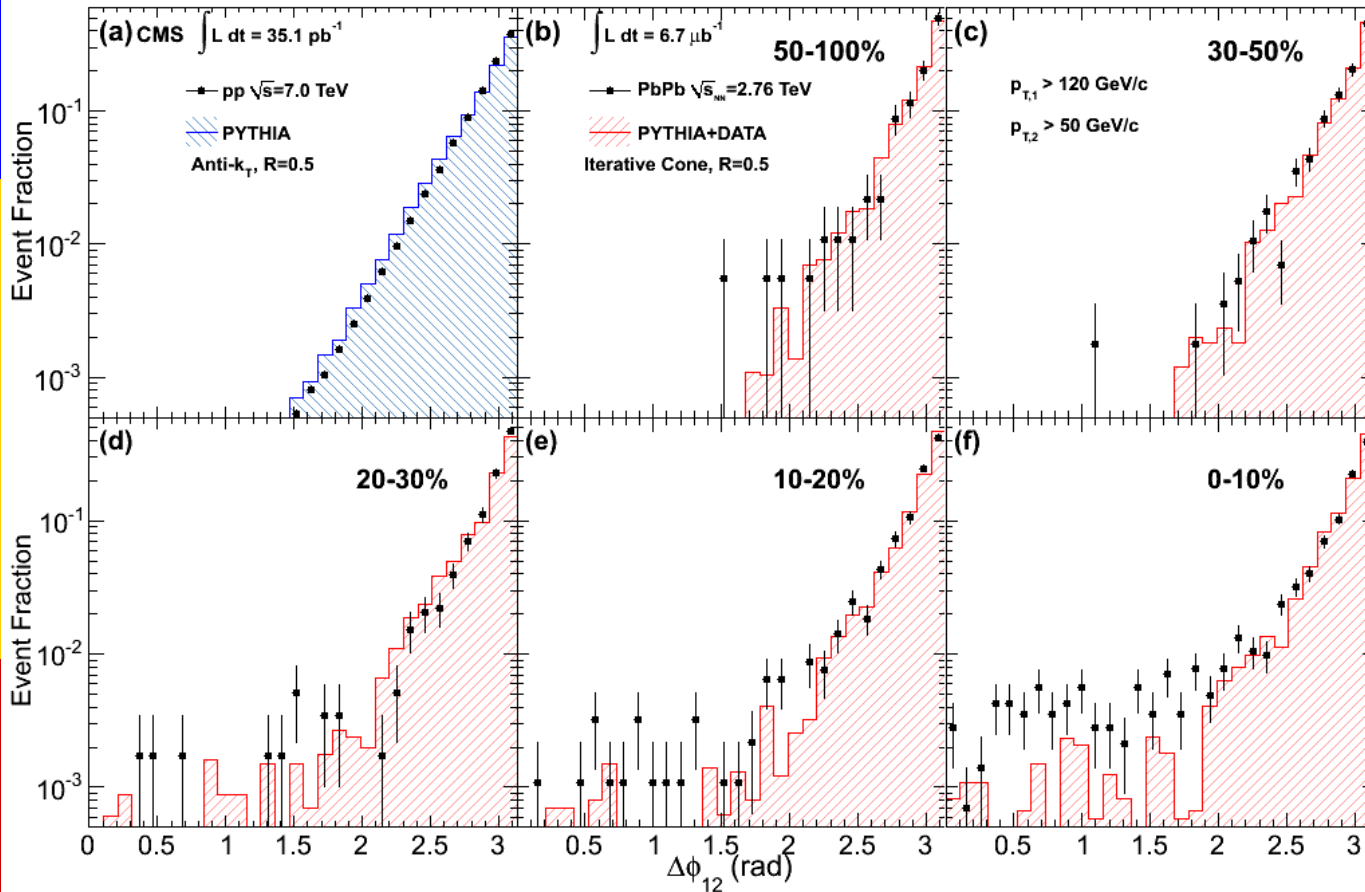
$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}} \quad \begin{array}{l} E_{T1} > 100 \text{ GeV} \\ E_{T2} > 25 \text{ GeV} \end{array}$$

- with increasing centrality:
  - enhancement of asymmetric di-jets with respect to pp
  - & HIJING + PYTHIA simulation
- $\Delta E \sim 20 \text{ GeV}$
- Consistent with RHIC



# Where is it lost ?

No visible angular decorrelation in  $\Delta\phi$  wrt pp collisions!



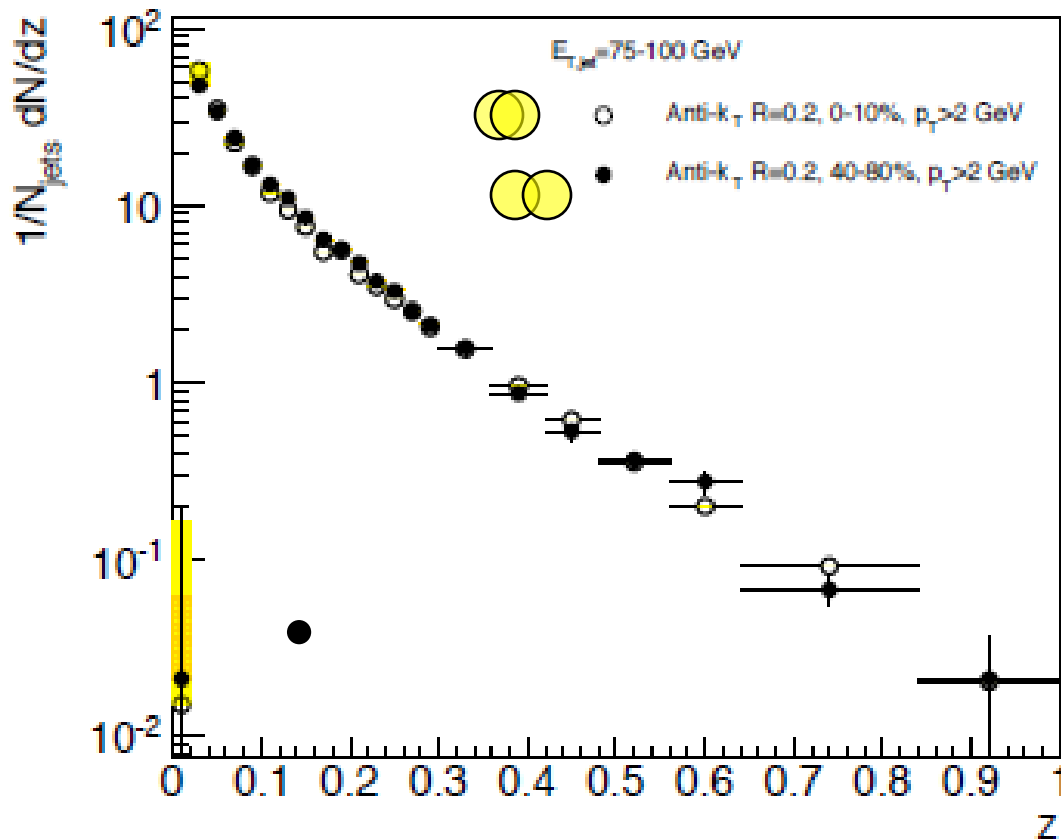
CMS: arXiv:1102.1957

→ large imbalance effect on jet energy, but very little effect on jet direction!

# How is it lost ?

## Jet fragmentation function

- distribution of the momenta of the fragments along the jet axis



Brian Cole – ATLAS (QM2011)

$$z = \frac{p_T^{hadron} \cdot \cos(\Delta R)}{E_T^{jet}}$$

- distribution is very similar in central and peripheral events
  - although quenching is very different...
- apparently no effect from quenching inside the jet cone...

# Jet “quenching”: what have we learned so far?

Large average dijet  $p_T$  imbalance

$p_T$  difference is found at low  $p_T$  far away from jet

Dijets remain back-to-back in azimuth

Partons fragment as in vacuum

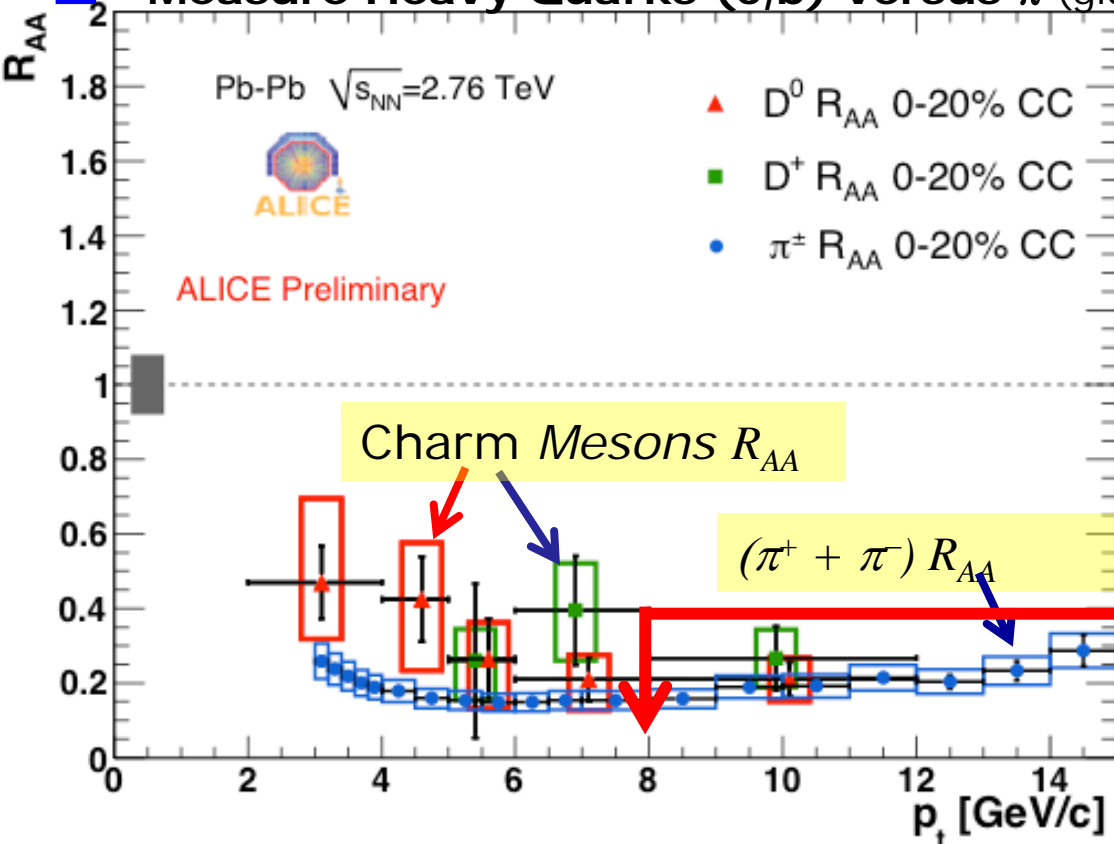
Experimental characterization of energy flow in dijet final states

*B. Wyslouch, CMS, EPIC2011*

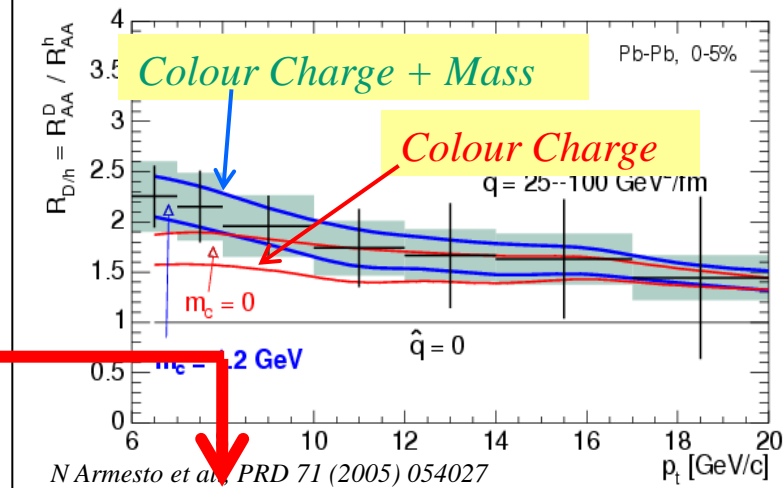
S. Mehdibadi

# Mass & Colour Charge Dependence

Measure Heavy Quarks (c,b) versus  $\pi$  (gluon fragmentation dominates  $\pi$  at LHC)



*D-Meson suppression /  $\pi$  suppression*



- $R_{AA}$  prompt charm  $\approx R_{AA}$  pions for  $p_T > 5-6$  GeV  
expected difference factor 2 @ 8 GeV
- $R_{AA}$  charm  $> R_{AA}$   $\pi$  for  $p_T < 5$  GeV ?

jet quenching

$$\Delta E \sim f(m) \times c_q \times \hat{q} \times L^2 \times f(E)$$

Needs better statistics & quantitative comparison with other models

# Summary

□ Journey started ~35 years ago: QGP was presumed

□ QGP observed at LHC

Our basic picture then is that matter at densities higher than nuclear consists of a quark soup.

The quarks become free at sufficiently high density. A specific realization is an asymptotically

■ **Collective flow**

□ **Strongly interacting liquid – hydro works**

□ Estimate of viscosity ?

■ **Jet quenching**

□ **Back to back jets strongly suppressed**

□ Dynamic of quenching ?

■ Quarkonia dissolution versus recombination ?

■ Heavy flavour: Are heavy quarks really suppressed as much as light quarks and gluons ?

□ **HI SM Model describing simultaneously all observables !**

J. C. Collins and M. J. Perry

*Mathematics and Theoretical Physics, University of Cambridge,  
Cambridge CB3 9EW, England*

(Received 6 January 1975)

**The journey to terra incognita of heavy ions continues!**

# Back up

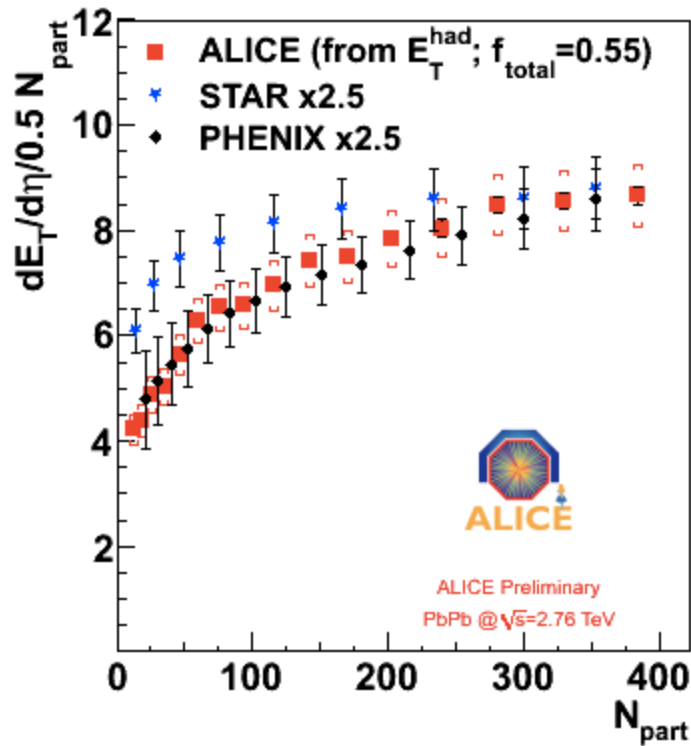




# Energy density

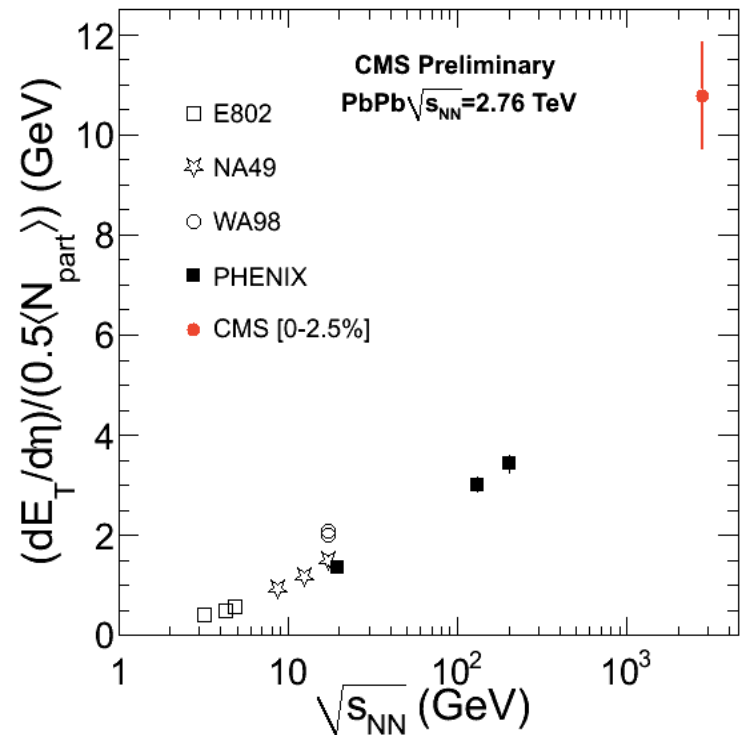
Transverse energy density per participant pair: 2.5 x RHIC

Consistent with 20% increase in  $\langle p_t \rangle$



J.Harris, ALICE, EPIC2011

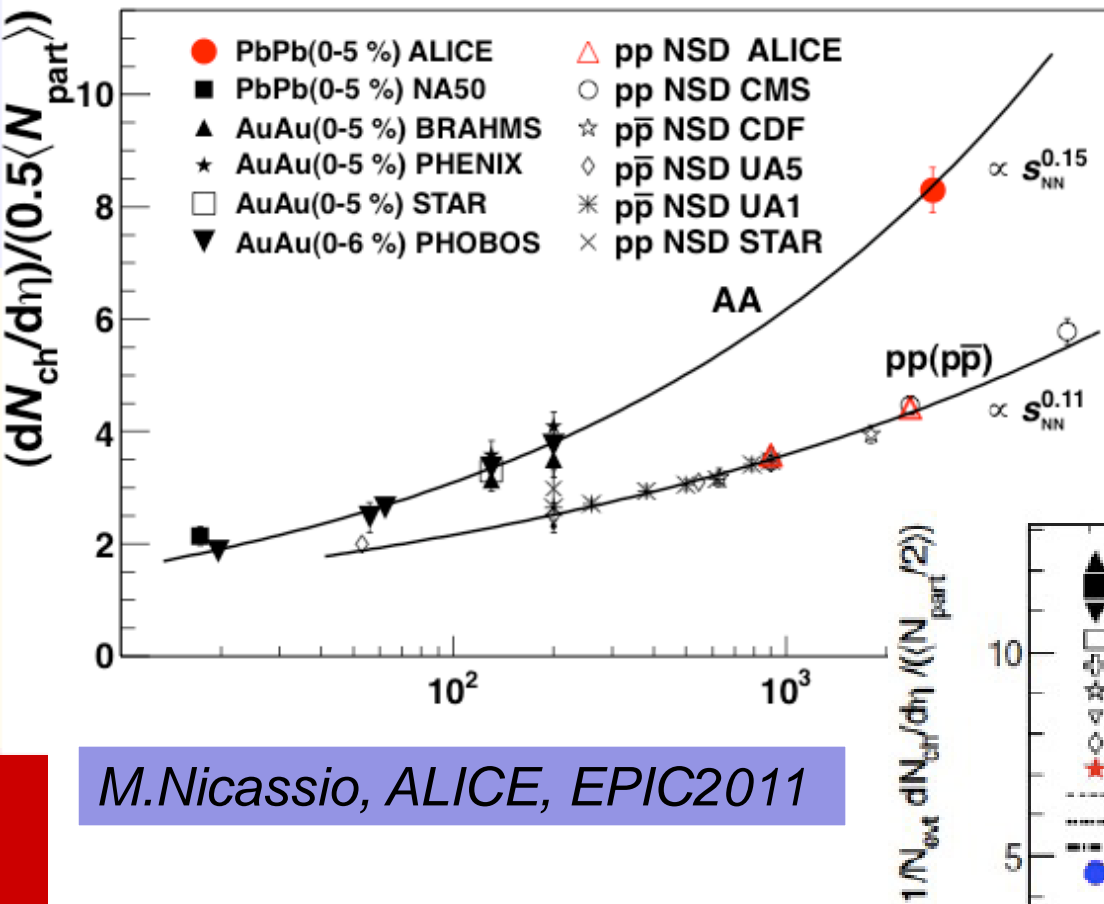
B.Wyslouch, CMS, EPIC2011



Bjorken energy density x time: 2.8  
for 5% of most central collisions

$$\epsilon = \frac{1}{\pi R^2 \tau} \frac{dE_T}{dy} \quad R = 1.12 A^{1/3} \text{ fm}$$

# Energy dependence

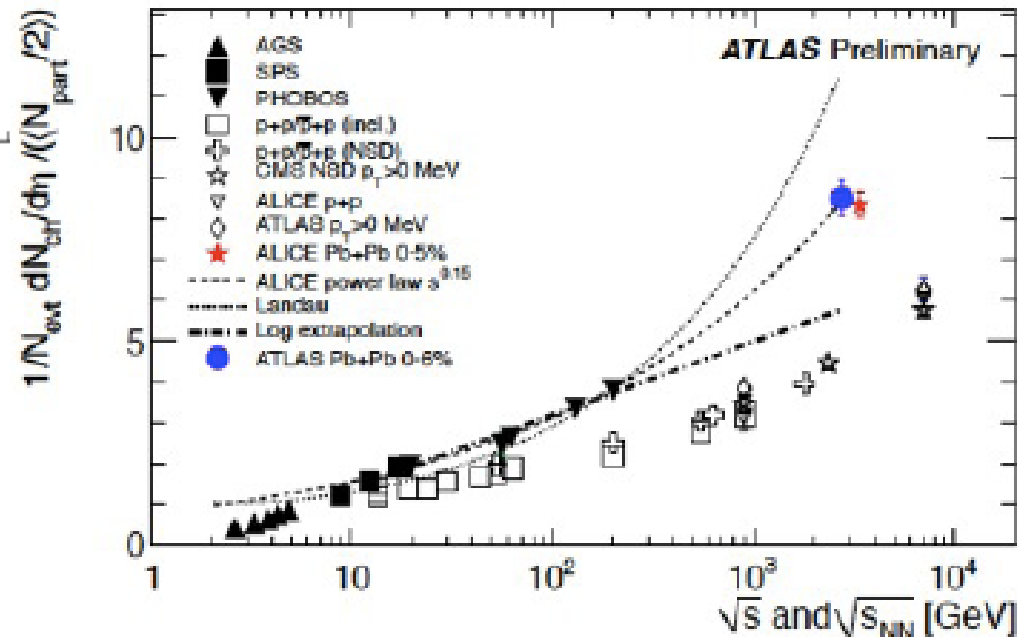


- growing faster than pp
- 2.2 compared to RHIC
- 1.9 compared to pp

S.White, ATLAS, EPIC2011

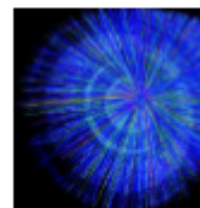
M.Nicassio, ALICE, EPIC2011

Agreement among experiments





# Charged Particle Multiplicity Density at mid-Rapidity

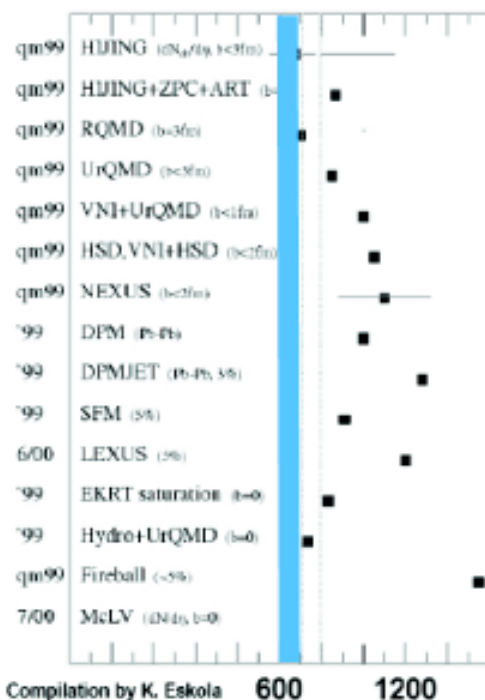


$$\left. \frac{dN_{ch}^{AA}}{d\eta} \right|_{\eta=0}$$

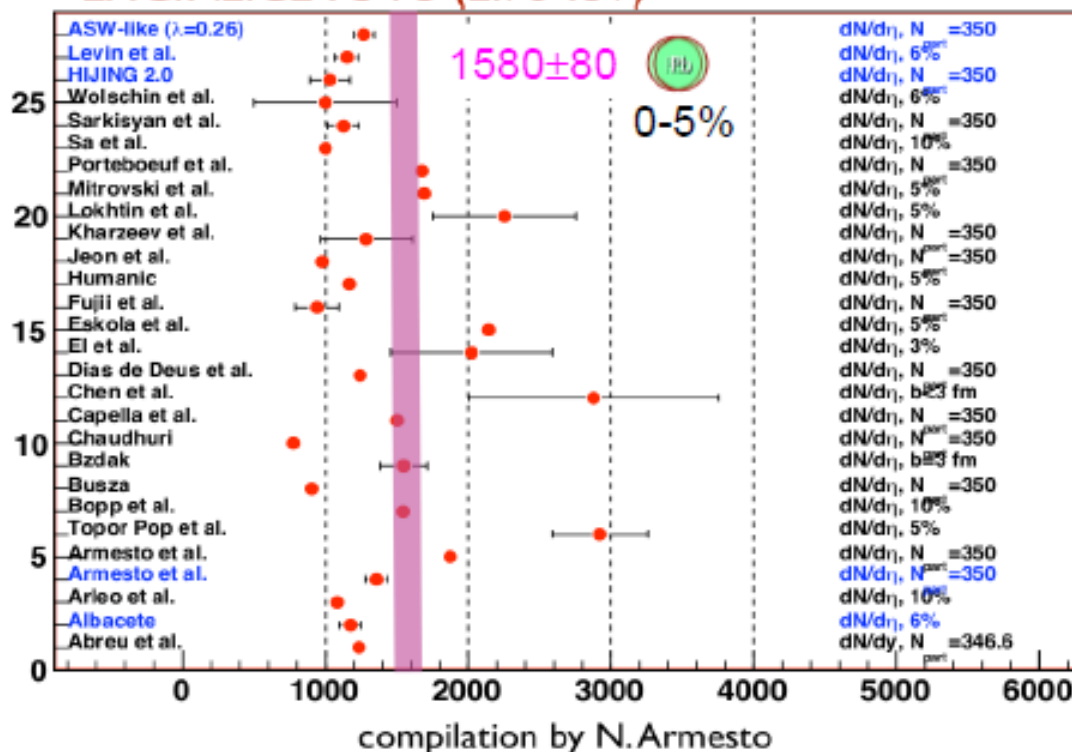
At RHIC the low charge particle multiplicity was a surprise and ruled out most of the models ...

No wonder, predictions for LHC were a little bit on the low side ...

RHIC: PHOBOS Au-Au (0.2 TeV)



LHC: ALICE Pb-Pb (2.76 TeV) Phys. Rev. Lett. 106, 032301 (2011)



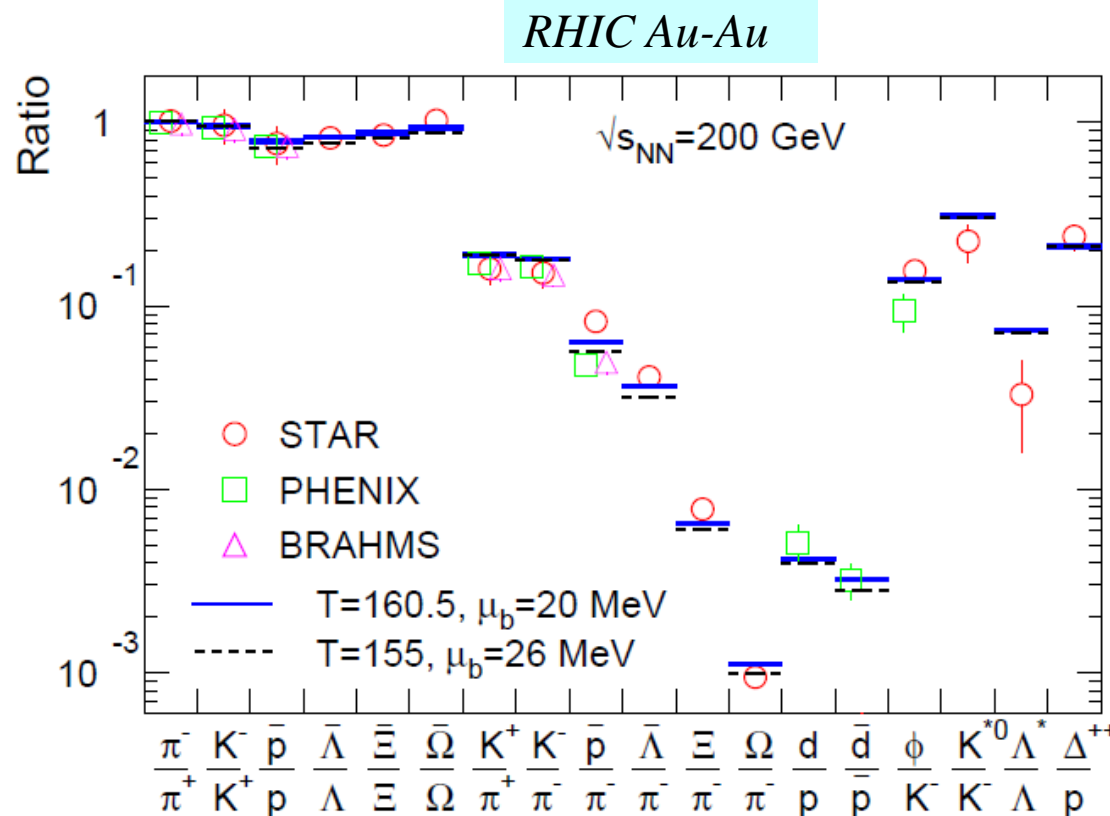
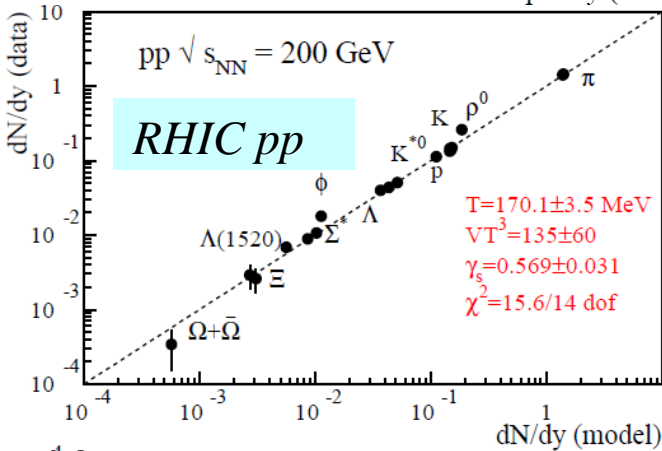
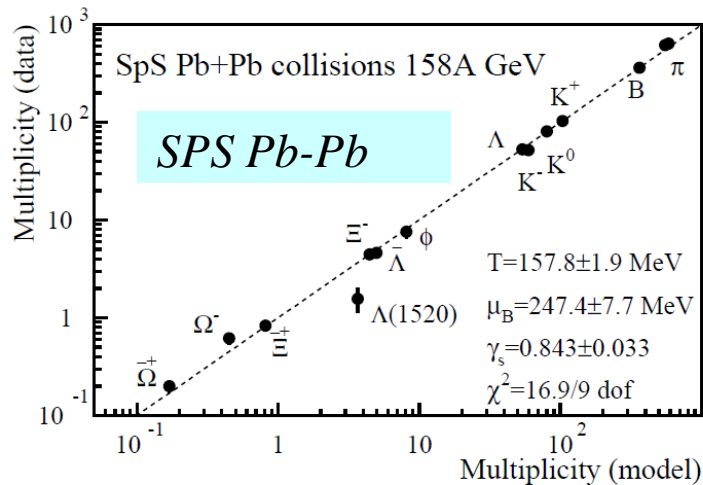
# Particle ratios

In general well described (~10%) by statistical (thermal) model

- $P(m) \sim e^{-(m/T)}$
- T Temperature
- $\mu_b$  Baryo-chemical potential
- $\gamma_s$  Strangeness suppression

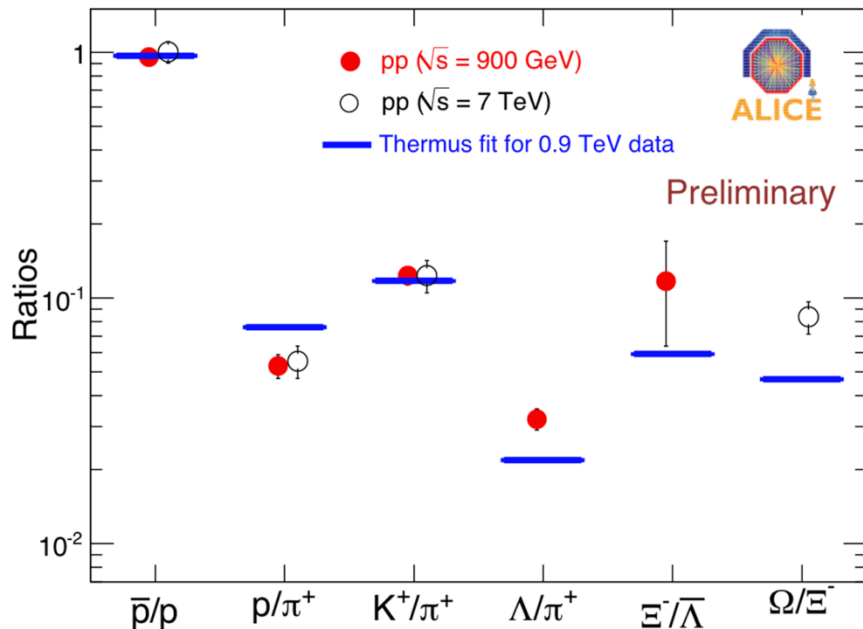
$$N_i/V = \frac{g_i}{(2\pi)^3} \gamma_s^{S_i} \int \frac{1}{\exp\left(\frac{E_i - \mu_B B_i - \mu_S S_i}{T_{\text{chem}}}\right) \pm 1} d^3 p,$$

$T_{ch}$ : 160-170 MeV     $\gamma_s$ : 0.9-1 (AA), 0.5-0.6 (pp)



# Thermal model at LHC

*pp: 900 GeV & 7 TeV*



**pp:**

- Thermal fit rather poor

**Pb-Pb:**

-  $K/\pi$  grows slightly from pp value

-  $p/\pi \approx$  like pp

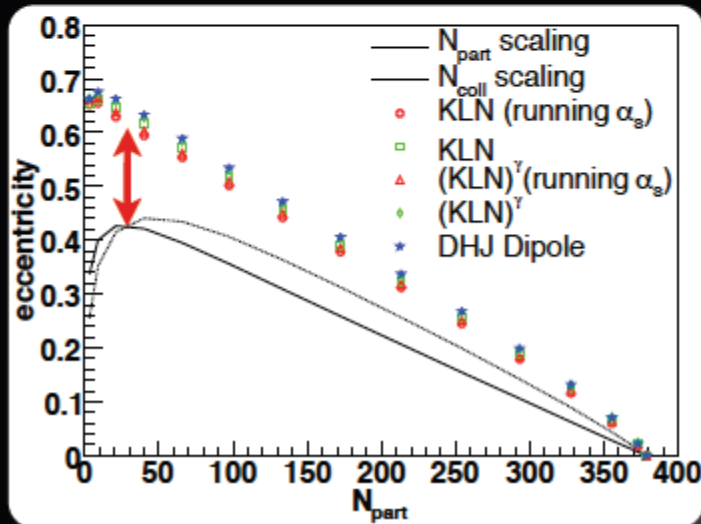
-  $p/\pi$  off by factor  $> 1.5$  from thermal predictions !

but very compatible with RHIC !!

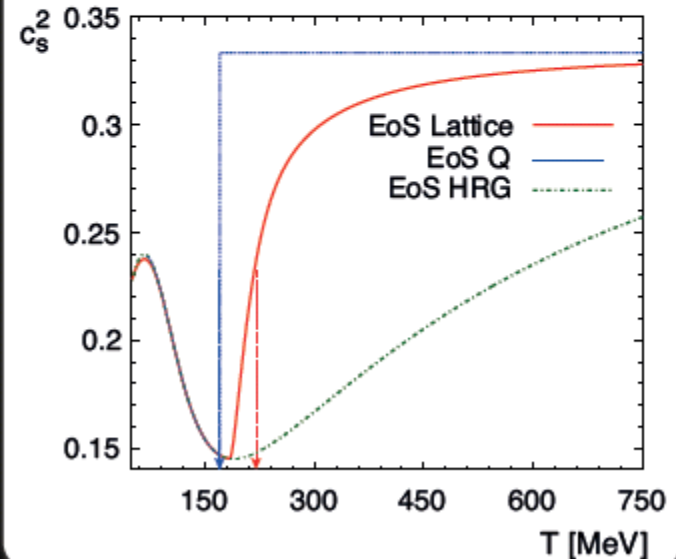
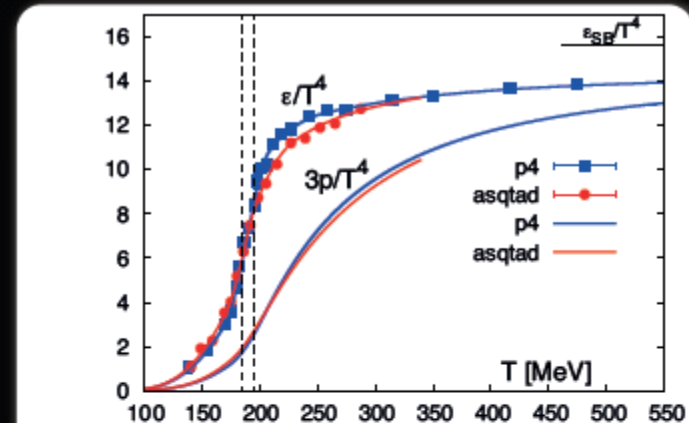
**Before we can conclude anything we need more particle species..**

# Physics of elliptic flow

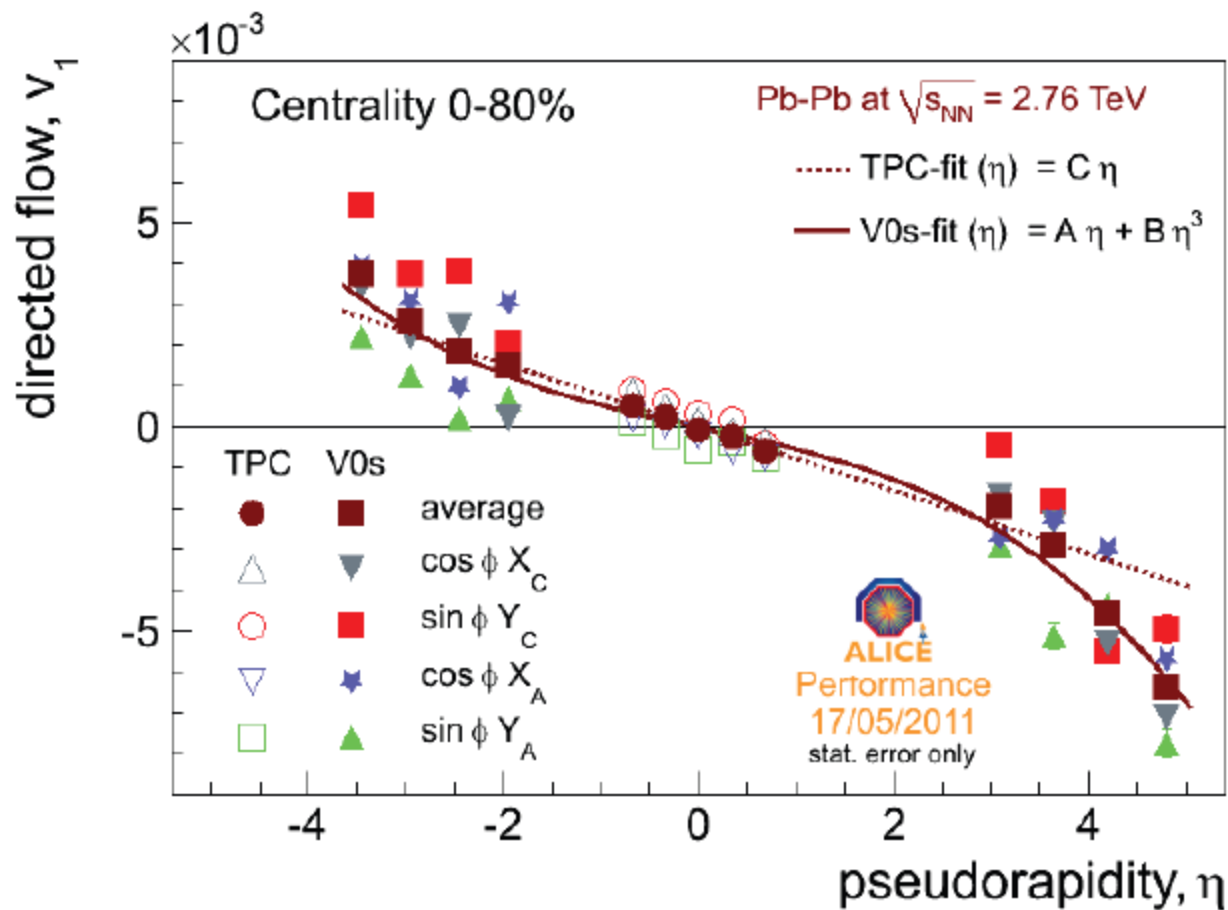
Elliptic flow  $v_2$  depends on fluid properties: the EoS via  $c_s^2 = \frac{\partial p}{\partial \epsilon}$ , shear viscosity over entropy ratio  $\eta/s$  but also on: initial conditions: particular initial spatial eccentricity  $\epsilon_2$



H.-J. Drescher et al.,  
Phys.Rev.C74:044905,2006

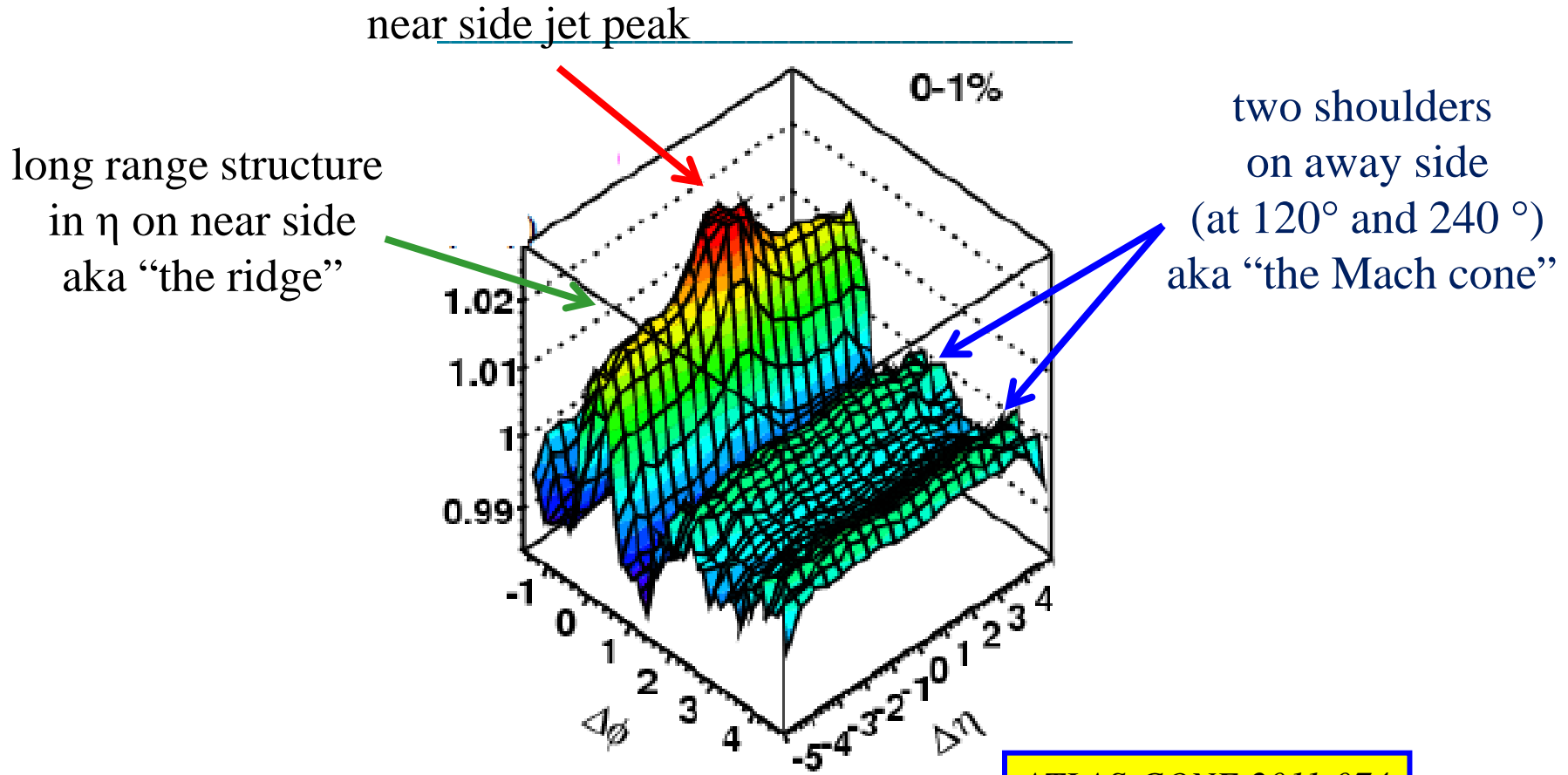


# Direct flow



G. Eyyubova QM2011

# Structures in $(\Delta\eta, \Delta\phi)$





# Fourier analysis at large $\Delta\eta$ , moderate $p_T$

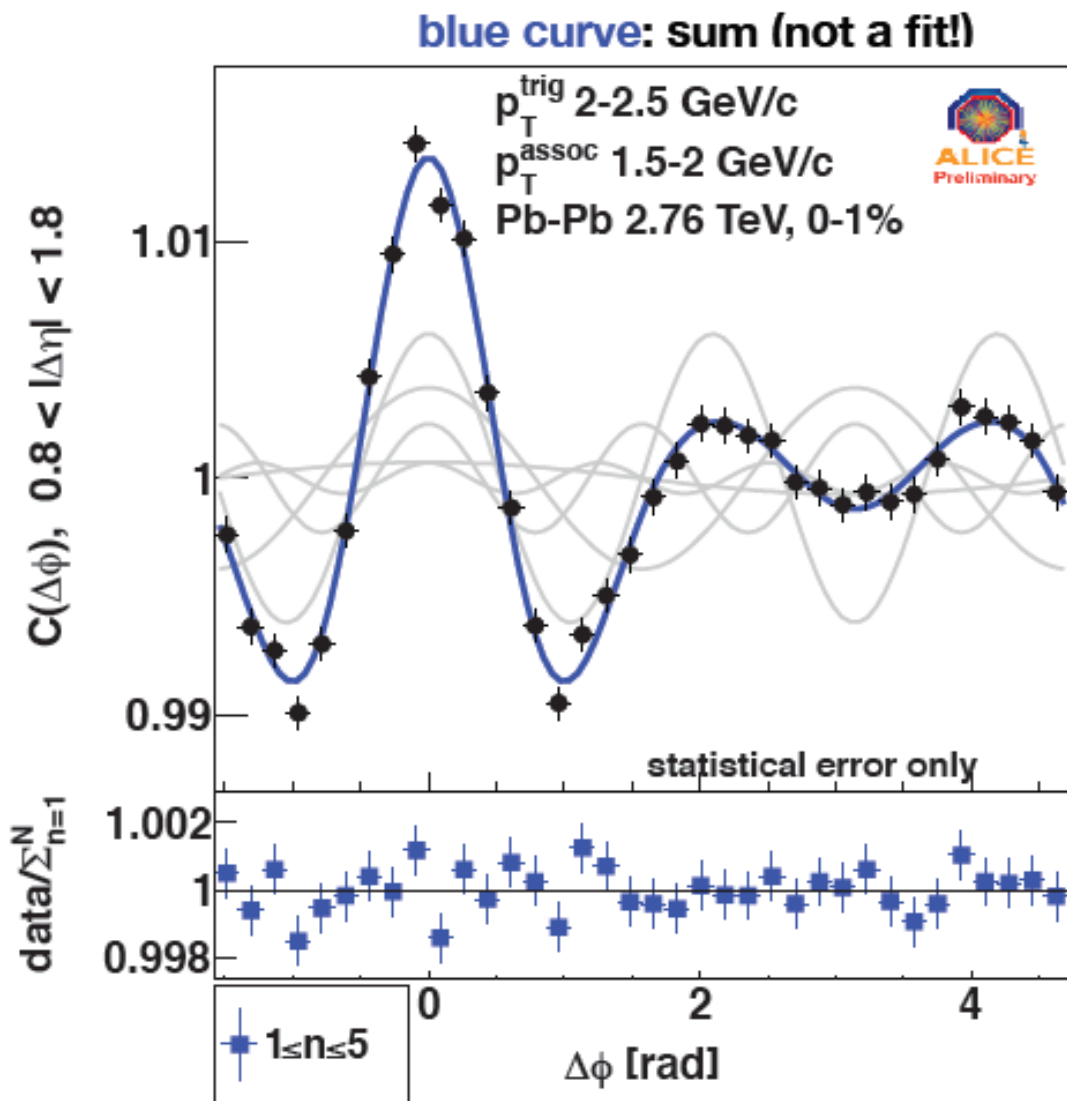
Near side jet excluded by  
 $|\Delta\eta| > 0.8$  gap  
 Ridge at  $\Delta\phi=0$  remains

2-particle Fourier coeffs.  
 Extract directly from  $C(\Delta\phi)$ :

$$\langle \cos n\Delta\phi \rangle = \frac{\int d\Delta\phi C(\Delta\phi) \cos n\Delta\phi}{\int d\Delta\phi C(\Delta\phi)}$$

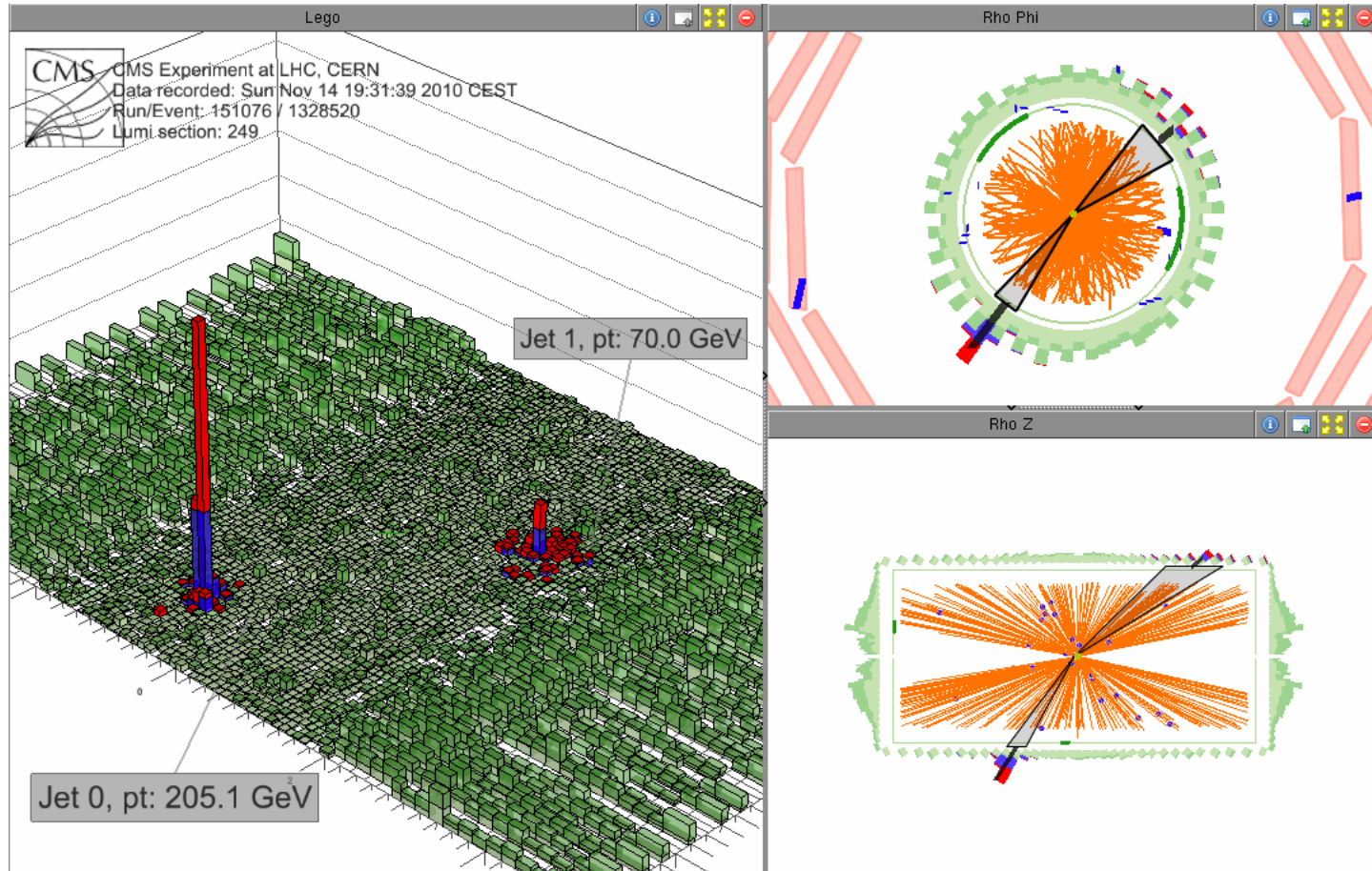
Here, the first 5 moments  
 describe shape at per-mille  
 level.

Andrew Adare – ALICE



# Di-jet imbalance

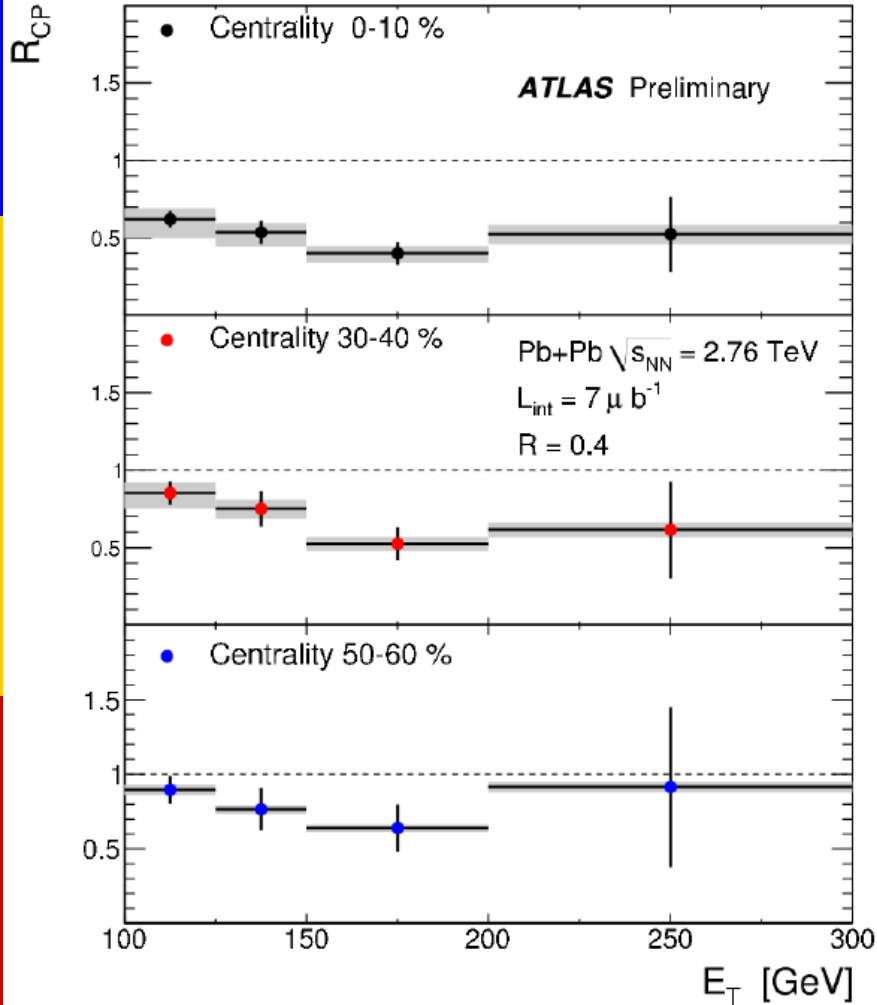
- Pb-Pb events with large di-jet imbalance observed at the LHC



→ recoiling jet strongly quenched!

*CMS: arXiv:1102.1957*

# Jet nuclear modification factor



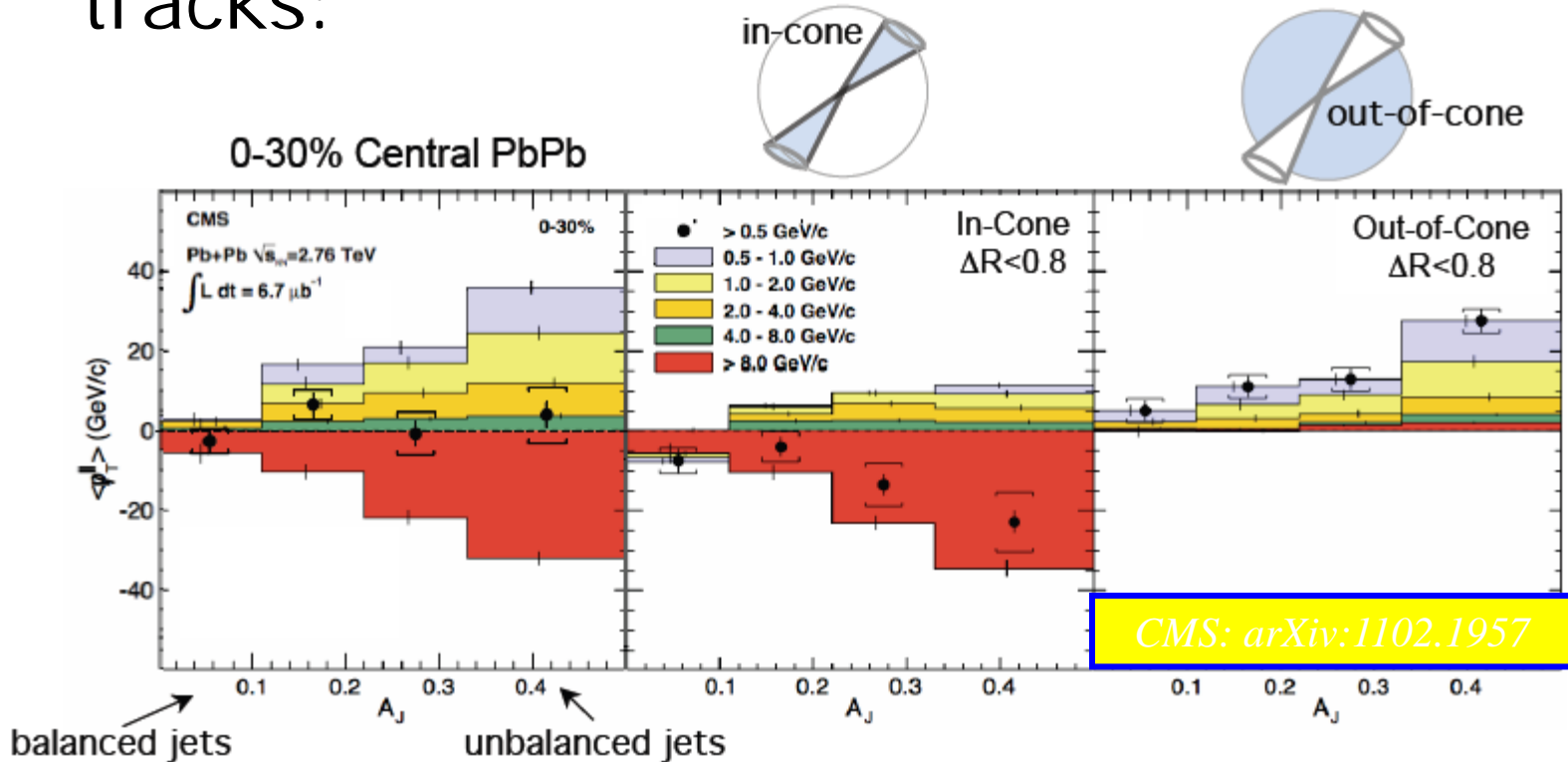
$$R_{CP} = \frac{\langle N_{bin} \rangle_{Central} \text{Yield}_{Central}}{\langle N_{bin} \rangle_{Peripheral} \text{Yield}_{Peripheral}}$$

- substantial suppression of jet production
  - in central Pb-Pb wrt binary-scaled peripheral
- out to very large jet energies!

Brian Cole – ATLAS (QM2011)

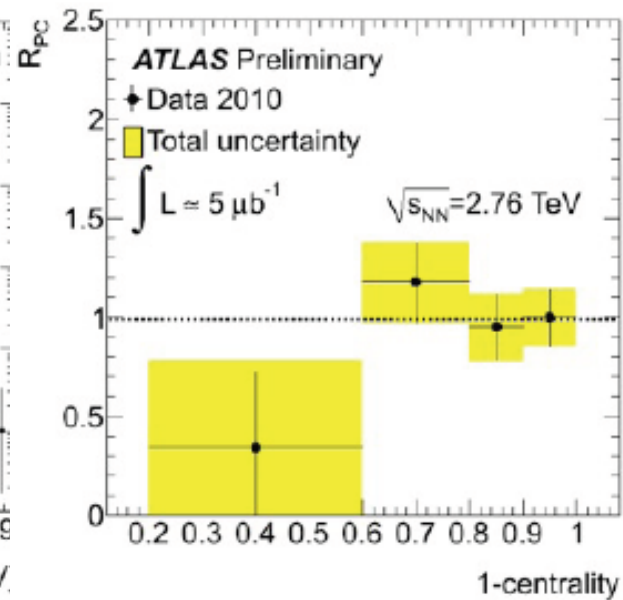
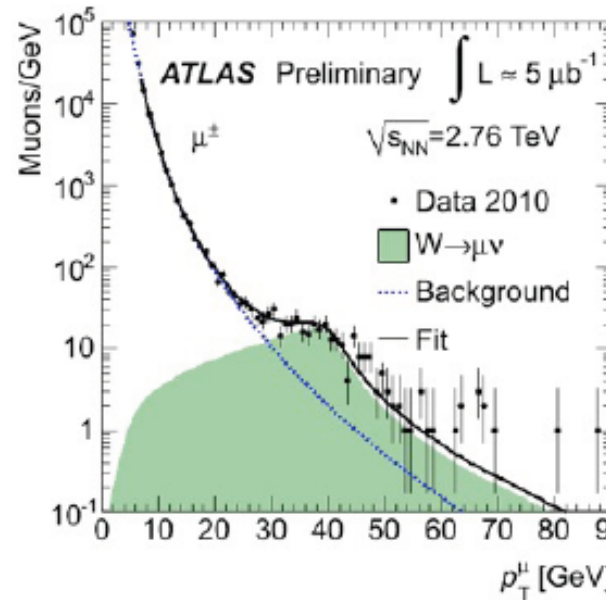
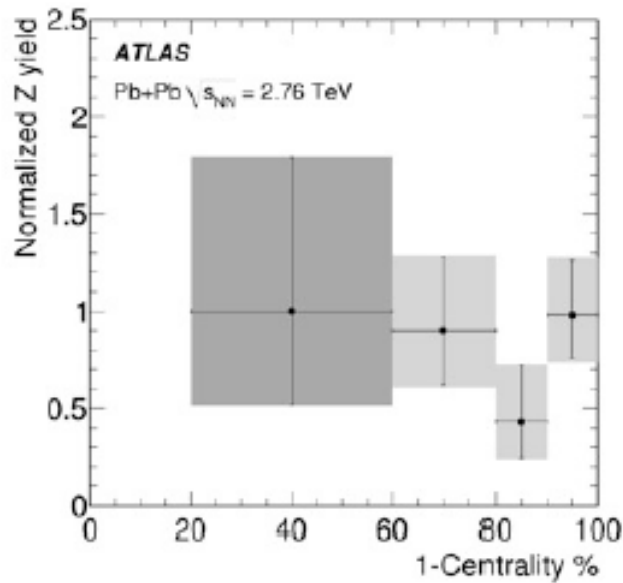
# Where does the energy end up?

- nice analysis by CMS using reconstructed tracks:



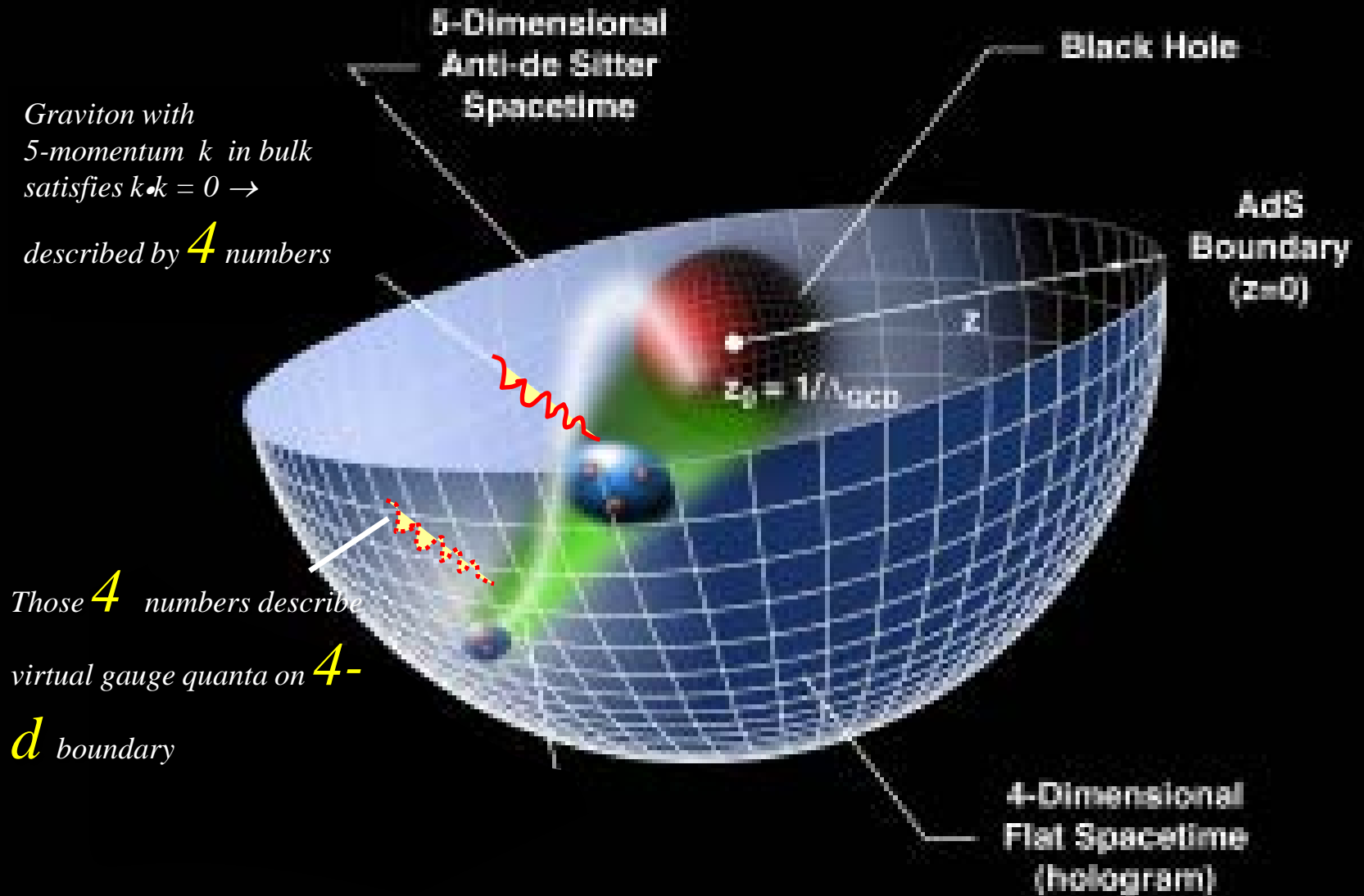
→ momentum difference is balanced by low momentum particles outside of the jet cone

# Z and W from ATLAS



*S. White, ATLAS, EPIC 2011*

# AdS / CFT in a Picture



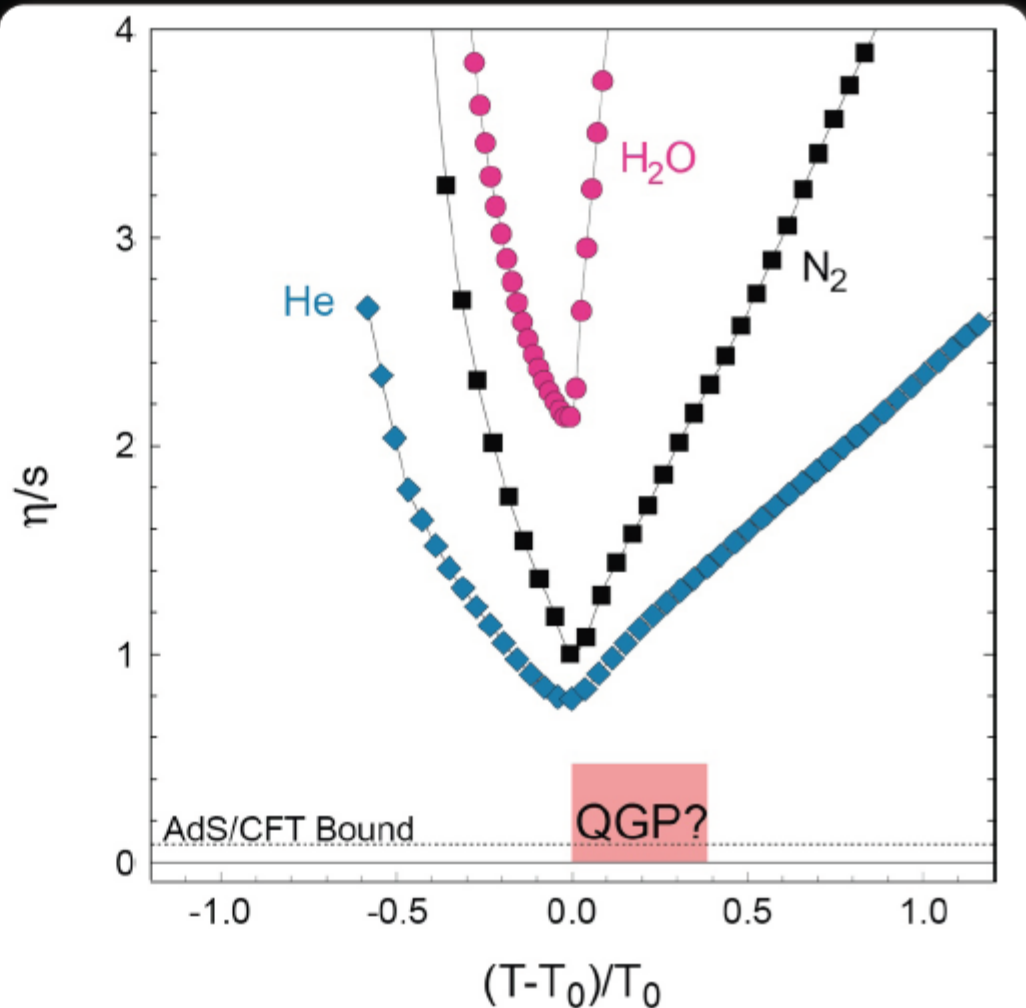
( Adopted from S. Brodsky figure )

# The Perfect Liquid?

model calculations suggest that the RHIC  $v_2$  results are close to the ideal hydrodynamical limit.

these calculations place an upper limit on  $\eta/s$  which is smaller than  $\sim 4 \times \text{AdS/CFT}$  bound

main uncertainties on  $\eta/s$  due to uncertainties in the initial conditions and the unknown dependence of  $\eta/s$  versus temperature



Based on R. Lacey et al., Phys.Rev.Lett.98:092301,2007.



# Heavy Ion Program

