# Studying the $\chi_b$ states with ATLAS Birmingham HEP Seminar

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- Introduction to the bottomonium system and some "History"
- Brief review of  $b\bar{b}$  spectroscopy
- $\chi_b$  studies at ATLAS
- $\chi_b$  analyses from DØ and LHCb
- Further opportunities at the LHC

#### Introduction: What are the $\chi_b$ states?

The  $\chi_b$  represent the spin triplet (S = 1) P-wave (L = 1) states of the bottomonium ( $b\bar{b}$ ) spectrum.

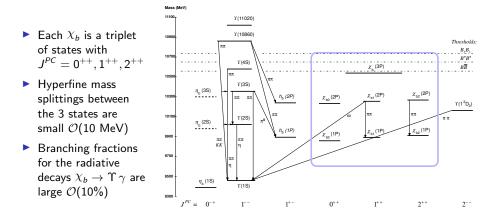


Figure: J. Beringer et al. (Particle Data Group), Phys. Rev. D86, 010001 (2012)

The radiative decays  $\chi_b \rightarrow \Upsilon(nS)\gamma$  represent the most experimentally clean channels to reconstruct  $\chi_b$  at the LHC:

- ▶ The radiative  $\chi_b$  decays benefit from large branching fractions  $\checkmark$
- $\mathcal{B}(\Upsilon(nS) \to \mu^+ \mu^-)$  is 1-2% and a di-muon signature is clean  $\checkmark$
- ▶ Photons are very soft X (< 1 GeV in X<sub>b</sub> rest frame)

PDG Summary (masses rounded to nearest 1 MeV,  $E_{\gamma}$  in rest frame of  $\chi_b$ )

|                 | Mass [MeV] | $E_{\gamma}$ [MeV] | $\mathcal{B}(\Upsilon(1S)\gamma)$ | $E_{\gamma}$ [MeV] | $\mathcal{B}(\Upsilon(2S)\gamma)$ |
|-----------------|------------|--------------------|-----------------------------------|--------------------|-----------------------------------|
| $\chi_{b0}(1P)$ | 9859       | 391                | <6%                               | -                  | -                                 |
| $\chi_{b1}(1P)$ | 9893       | 423                | 35%                               | -                  | -                                 |
| $\chi_{b2}(1P)$ | 9912       | 442                | 22%                               | -                  | -                                 |
| $\chi_{b0}(2P)$ | 10233      | 743                | 1%                                | 207                | 5%                                |
| $\chi_{b1}(2P)$ | 10255      | 764                | 9%                                | 230                | 21%                               |
| $\chi_{b2}(2P)$ | 10269      | 777                | 7%                                | 242                | 16%                               |

# Predictions from QCD inspired potential models:

- 1. Phys. Rev. D 36 3401 (1987)
- 2. Phys. Rev. D 38 279 (1988)
- 3. Eur. Phys. J. C. 4 107 (1998)
- Just below the  $B\overline{B}$  threshold (10.558 GeV)
- Narrow, Γ < 1 MeV if...</p>
- ... $\Gamma(\chi_b(3P) \to \Upsilon(1,2,3S)\gamma)/\Gamma_{Tot.}$  is large (expected to be so)
- $\succ \ \Gamma(\Upsilon(3S)\gamma) > \Gamma(\Upsilon(2S)\gamma) > \Gamma(\Upsilon(1S)\gamma)$

| Publication | 3 <sup>3</sup> <i>P</i> c.o.g | $\Delta 3^3 P_0$ | $\Delta 3^3 P_1$ | $\Delta 3^3 P_2$ |
|-------------|-------------------------------|------------------|------------------|------------------|
| 2.          | 10.520 GeV                    | -19 MeV          | -4 MeV           | +6 MeV           |
| 3.          | 10.525 GeV                    | -22 MeV          | -4.9 MeV         | +7.3 MeV         |

#### History: 1974 - Lederman et al. propose E288

NAL PROPOSAL # 288

Scientific Spokesman:

L. M. Lederman Physics Department Columbia University New York, New York 10027

PTS/Off-net: 212 - 460-0100 280-1754



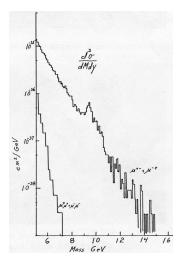
A Study of Di-Lepton Production in Proton Collisions at NAL

J. A. Appel, M. H. Bourquin, D. C. Hom, L. M. Lederman, J. P. Repellin, H. D. Snyder, J. K. Yoh (Columbia University); B. C. Brown, P. Limon, T. Yamanouchi (NAL).

(Formerly #70 Phase III)

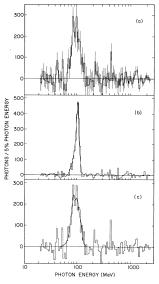
- Observe and measure the spectrum of virtual photons emitted in p-nucleon collisions via the mass distribution of e<sup>+</sup>e<sup>-</sup> pairs: p + p + e<sup>+</sup>e<sup>-</sup> + anything. (1) Study characteristics, e.g. parity violation, p<sub>1</sub> behavior.
- Search for structures in the above spectrum, publish these and become famous, e.g. W\*, B\*.
- Extend the Experiment \$70 study of single leptons in the double arm arrangement, i.e. W<sup>±</sup> etc. Publish these and become famous.

## History: 1977 - The CFS Collaboration discover the $\Upsilon$ with E288





"An experimental group at the Fermi National Accelerator Laboratory announced recently that it has discovered a new particle. The new particle has a mass of 9.5 GeV..."



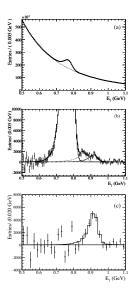


Phys. Rev. Lett. 49, 1612 and 1616 (1982)

- ► First saw evidence of  $\Upsilon(3S) \rightarrow \chi_b(2P)\gamma$  (figure)
- ▶ Then observed the  $\chi_b(2P) \rightarrow \Upsilon(1,2S)\gamma$

#### Phys. Rev. Lett. 51, 160 (1983)

► The next year,  $\Upsilon(2S) \rightarrow \chi_b(1P)\gamma$  and  $\chi_b(1P) \rightarrow \Upsilon(1S)\gamma$  were both observed

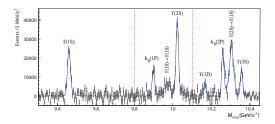


Phys. Rev. Lett. 101, 071801 (2008)

• Babar observed  $\Upsilon(3S) \rightarrow \eta_b(1S)\gamma$  (left)

Phys. Rev. Lett. 108, 032001 (2012)

▶ Belle observed  $\Upsilon(5S) \rightarrow h_b(1, 2P)\pi^+\pi^+$ (below)



All discoveries in the bottomonium system since the  $\Upsilon$  made by  $e^+e^-$  experiments!

#### Phys. Rev. Lett. 108, 122001 (2012) (arXiv:1110.2251)

Belle recently reported the observation of two narrow structures in  $\pi^{\pm}\Upsilon(nS)$  (n = 1, 2, 3) and  $\pi^{\pm}h_b(mS)$  (m = 1,2) pairs produced in association with a single charged pion in  $\Upsilon(5S)$  decays!

Quarkonium physics can still surprise us!

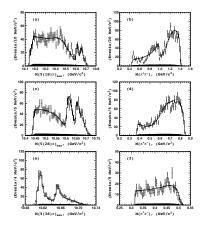
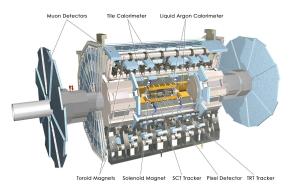


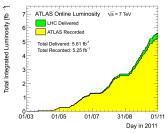
FIG. 2: Comparison of fit results (open histogram) with experimental data (points with error bars) for events in the  $\Upsilon(1S)$  (a,b),  $\Upsilon(2S)$  (c,d), and  $\Upsilon(3S)$  (e,f) signal regions. The hatched histogram shows the background component.

## The ATLAS Detector at the LHC

The ATLAS detector is a general purpose particle physics detector designed to study physics at the TeV scale:



ATLAS has a diverse physics programme including Higgs Searches, SUSY + Exotics Searches, SM Physics, Heavy Flavour Physics and more! The LHC and ATLAS performed very well throughout 2011:



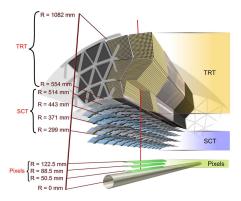
ATLAS collected **over** 5 fb<sup>-1</sup> of data during the 2011 LHC run at  $\sqrt{s} = 7$  TeV

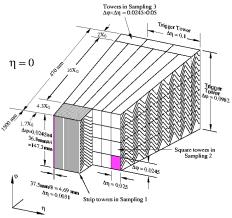
Also, over 14 fb<sup>-1</sup> collected to date at  $\sqrt{s} = 8$  TeV during 2012!

#### Detector Components I

#### Inner Detector (ID) ( $|\eta| < 2.5$ )

 Silicon Pixels and Strips (SCT) with Transition Radiation Tracker (TRT)





#### Liquid Argon EM Calorimeter ( $|\eta| < 3.2$ )

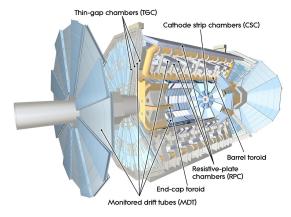
 Highly granular and longitudinally segmented in 3-4 layers

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|-------------|---|---------|
|-------------|---|---------|

## Detector Components II

#### Muon Spectrometer (MS) ( $|\eta| < 2.7$ )

Toroid Magnet, 4 detector technologies, dedicated tracking and trigger chambers



- Barrel: MDT (Tracking) and RPC (Trigger)
- Endcaps: MDT + CSC (Tracking) and TGC (Trigger)

# Observation of a new $\chi_b$ state in radiative transitions to $\Upsilon(1S)$ and $\Upsilon(2S)$ at ATLAS Phys. Rev. Lett. 108, 152001 (2012) (arXiv:1112.5154 [hep-ex])

Radiative  $\chi_b$  decays are studied with two simultaneous analyses which exploit different reconstruction methods and detectors:

- Photons reconstructed using the EM calorimeter (denoted unconverted)
- ▶  $\gamma \rightarrow e^+e^-$  conversions reconstructed with the Inner Detector (denoted converted)
- Both share a common  $\Upsilon \to \mu^+ \mu^-$  selection

The two reconstruction methods have their own advantages and disadvantages. In particular, the minimum  $p_{T}(\gamma)$  threshold (2.5 GeV and 1.0 GeV respectively) determines which radiative decays can be reconstructed:

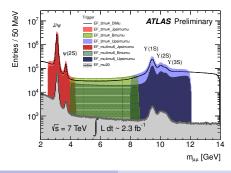
- ► The unconverted photon analysis is capable of reconstructing  $\chi_b \rightarrow \Upsilon(1S) \gamma$  decays alone
- The converted photon analysis is capable of reconstructing both *χ<sub>b</sub>* → Υ(1*S*) γ and *χ<sub>b</sub>* → Υ(2*S*) γ decays

## Data Sample and Trigger Selection

The analysis uses 4.4 fb<sup>-1</sup> of *pp* collision data at  $\sqrt{s} = 7$  TeV recorded throughout the 2011 LHC run:

Trigger Strategy:

- Events containing radiative  $\chi_b$  decays are triggered by the di-muon decay  $\Upsilon \rightarrow \mu^+ \mu^-$  (the photons are too soft to trigger the event)
- The trigger records events which contain di-muon pairs or single high  $p_T$  muons
- ► The majority of events are selected by dedicated \(\U03c4 → \mu^+ \mu^-\) di-muon triggers (blue shaded histograms)

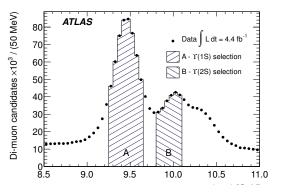


## Common $\Upsilon$ Selection

Selection of  $\Upsilon(1,2S) \rightarrow \mu^+ \mu^-$  candidates is common to both the unconverted and converted photon analyses:

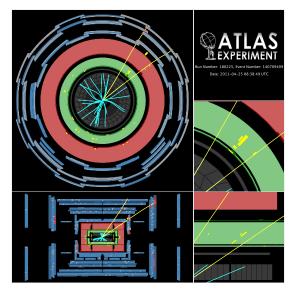
#### Muon Selection

- ▶ p<sub>T</sub>(µ<sup>±</sup>) > 4.0 GeV
- |η(μ<sup>±</sup>)| < 2.3
  </p>
- Reconstructed from track in ID combined with MS track
- $\Upsilon \to \mu^+ \mu^-$  Selection
  - Oppositely charged di-muon pair
  - $\mu^+\mu^-$  common vertex fit  $\chi^2/N_{D.o.F} < 20$
  - $p_T(\mu^+\mu^-) > 12 \text{ GeV}$
  - Rapidity  $|y(\mu^+\mu^-)| < 2.0$
  - Both muons associated to same primary pp interaction



 $\frac{\Upsilon \to \mu^+ \mu^- \text{ invariant mass selection}}{\blacktriangleright \text{ A - } \Upsilon(1S): 9.25 < m(\mu^+ \mu^-) < 9.65 \text{ GeV}}$  $\blacktriangleright \text{ B - } \Upsilon(2S): 9.80 < m(\mu^+ \mu^-) < 10.10 \text{ GeV}$ A. Chisholm Studying the  $\chi_b$  states with ATLAS 16 / 45

## **Unconverted** Photon Analysis



An event containing a candidate  $\chi_b\to\Upsilon\gamma$  decay in which the photon is unconverted

## Unconverted Photon Selection

EM calorimeter energy deposits not matched to any track are considered as **unconverted** photon candidates:

- ► E<sub>T</sub>(γ) > 2.5 GeV
- ▶  $|\eta(\gamma)| < 2.37$  (Barrel-Endcap transition region  $1.37 < |\eta| < 1.52$  excluded)
- "Loose"<sup>†</sup> photon ID selection: Including limits on hadronic leakage and requirements on the EM shower shape (designed to reject backgrounds from narrow jets and π<sup>0</sup> decays)

#### **Unconverted** Photon Pointing Correction

- The polar angle of the photon 3-vector is corrected to point back to μ<sup>+</sup>μ<sup>-</sup> vertex
- ► Loose cut of \u03c0<sup>2</sup>/N<sub>D.o.F</sub> < 200 rejects photons not compatible with having originated from the \u03c0<sup>+</sup>\u03c0<sup>-</sup> vertex

 $\chi_b \to \Upsilon(1S) \gamma$  Selection

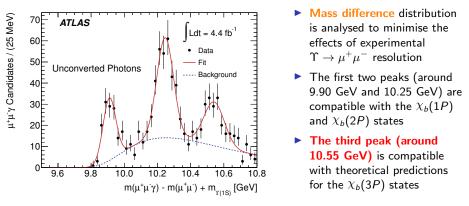


► Reconstructed  $\Upsilon(1S) \rightarrow \mu^+ \mu^-$  candidates are associated with corrected unconverted photons to form  $\chi_b$  candidates

<sup>†</sup>Described in detail in: Phys. Rev. D 83, 052005 (2011) (arXiv:1012.4389)

#### Unconverted Photon Result I

The resulting  $m(\mu^+\mu^-\gamma) - m(\mu^+\mu^-) + m_{\Upsilon(1S)}^{PDG}$  distribution exhibits three peaks:



- ► Final selection of p<sub>T</sub>(µ<sup>+</sup>µ<sup>-</sup>) > 20 GeV chosen to maximise X<sub>b</sub>(1P) and X<sub>b</sub>(2P) significance irrespective of effect on the third peak
- Statistical significance of third signal is greater than 6σ calculated from a likelihood ratio approach (including systematic variations)

An extended unbinned maximum likelihood fit is performed to the  $m(\mu^+\mu^-\gamma) - m(\mu^+\mu^-) + m_{\Upsilon(1S)}^{PDG}$  distribution to extract an estimate of the  $\chi_b(3P)$  mass barycentre:

Fit Model

- **Signal:** Single Gaussian for each  $\chi_b(nP)$  peak, each with a free mean value and width
- Background: Described by exp (A · (ΔM) + B · (ΔM)<sup>-2</sup>) where A and B are free parameters

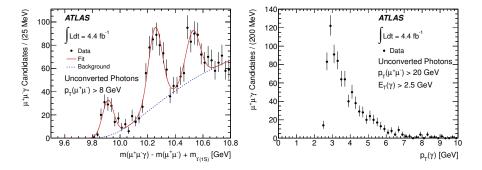
#### **Assigned Systematic Uncertainties**

- Unconverted photon energy scale uncertainty (estimated at ±2% of the ΔM position)
- Modelling of the background distribution (estimated from refitting with various alternative models)

|              | Fitted Mass (MeV)                             |
|--------------|---|
| $\chi_b(1P)$ | $9910\pm 6~(	ext{stat.})\pm 11~(	ext{syst.})$ |
| $\chi_b(2P)$ | $10246\pm5~(	ext{stat.})\pm18~(	ext{syst.})$  |
| $\chi_b(3P)$ | $10541\pm11$ (stat.) $\pm$ 30 (syst.)         |

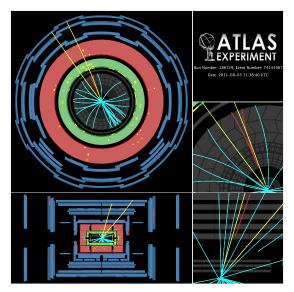
The statistical significance of third signal remains greater than  $6\sigma$  with each systematic variation

Left:  $m(\mu^+\mu^-\gamma) - m(\mu^+\mu^-) + m_{\Upsilon(1S)}^{PDG}$  distribution without a lower  $p_T(\mu^+\mu^-) > 8$  GeV cut.



**Right:** Unconverted photon  $p_T(\gamma)s$  distribution for  $\Upsilon(1S)\gamma$  candidates.

## **Converted** Photon Analysis

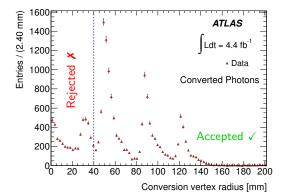


An event containing a candidate  $\chi_b o \Upsilon \gamma$  decay in which the photon has converted ( $\gamma o e^+ e^-$ )

## Converted Photon Selection I

Reconstructing photons from  $e^+e^-$  conversions in the Inner Detector (ID) offers improved resolution and access to softer photons:

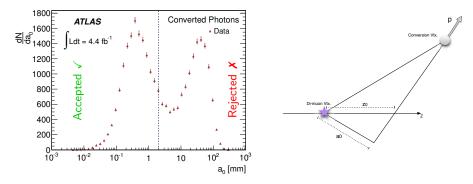
- Reconstructed from ID measurements *alone* (no EM cluster matching)
- Minimum track momentum p<sub>T</sub>(e<sup>±</sup>) > 500 MeV
- ▶  $p_T(\gamma) > 1$  GeV
- |η(γ)| < 2.3</p>
- Only two-track conversions are retained
- 4 silicon detector hits required for each electron track



- Candidate electron tracks must not already be selected as di-muon candidate tracks
- ▶ Radius of conversion vertex R > 40 mm to reduce background contamination

## Converted Photon Selection II

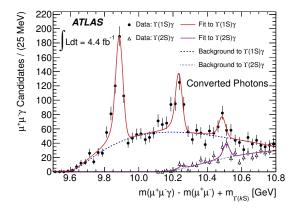
The 3D impact parameter of the converted photon with respect to the di-muon vertex,  $a_0$ , is a powerful variable which can be used to select photons associated with the di-muon vertex:



▶ a<sub>0</sub> < 2 mm is required to reject photon combinatorics not compatible with having originated from the di-muon vertex

• The  $\chi^2$  probability of the conversion vertex fit is required to be greater than 0.01

Both the  $\chi_b \to \Upsilon(1S) \gamma$  and  $\chi_b \to \Upsilon(2S) \gamma$  distributions are shown together:



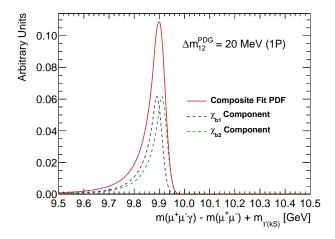
- Statistical significance of the third signal (around 10.5 GeV) is greater than 6σ calculated from a likelihood ratio approach (including systematic variations)
- Data points are not corrected for energy losses due to Bremsstrahlung (taken into account in fit)

Under the interpretation of the third signal as  $\chi_b(3P)$ , the experimental mass barycentre is measured from a simultaneous unbinned extended maximum likelihood fit to both the  $\Upsilon(1S)\gamma$  and  $\Upsilon(2S)\gamma$  mass distributions:

 The simultaneous fit allows a number of parameters to be shared across the two samples to help constrain the model, with additional constraints applied from the known masses (PDG)

Fit Model:

- As the J = 0 branching fraction is significantly smaller than for J = 1, 2 its contribution can be neglected
- The \u03c6<sub>b</sub>(nP) state is therefore modelled by two Crystal Ball (CB) functions to describe the low-mass Bremsstrahlung tail
- ▶ For n = 1, 2, the masses of the individual J=1,2 states are fixed to the known PDG values, and for n=3 the hyperfine splitting is fixed to the theoretically predicted value of 12 MeV
- ▶ The relative normalisations of the J=1 and J=2 components are fixed to be equal
- A free parameter  $\lambda$ , common to all the peaks, accounts for additional energy losses and appears in the form  $\overline{\Delta m} \cdot \lambda$
- The background is modelled by  $(\Delta m q_0)^{\alpha} \cdot \exp \{(\Delta m q_0) \cdot \beta\}$



- For demonstration,  $\sigma = 20$  MeV (i.e. 1P J = 1, 2 splitting)
- No knowledge of  $\sigma \cdot \mathcal{B}$  for any of the states
- Relative normalisation of \(\chi\_{b1}\) and \(\chi\_{b2}\) components is fixed to be equal\)

## Converted Photon Result II

#### Assigned Systematic Uncertainties:

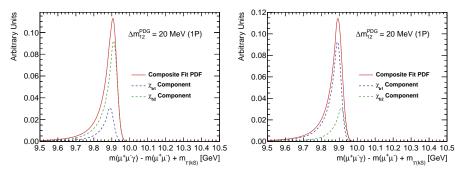
- ▶ Vary relative J = 1, 2 signal normalisation by  $\pm 0.25$  (or left free in fit):  $\pm 5$  MeV
- Alternative signal and background models: ±5 MeV
- Decoupled fits to the  $\Upsilon(1S)$  and  $\Upsilon(2S)$  distributions:  $\pm 5$  MeV
- Individually releasing constraints to the PDG values for the  $\chi_b(1P)$  and  $\chi_b(2P)$  masses:  $\pm 3$  MeV

#### Fit Result:

- Energy scale factor  $\lambda = 0.961 \pm 0.003$
- Experimental mass barycentre for  $\chi_b(3P)$  signal determined by fit to converted photon candidates **alone** is:

$$\overline{m}_3 = 10.530 \pm 0.005 \text{ (stat.)} \pm 0.009 \text{ (syst.)}$$
 GeV

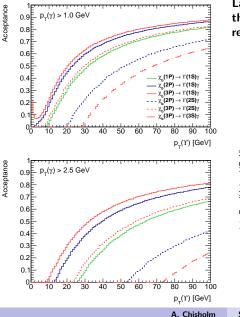
## Demonstration of J = 1, 2 Normalisation Systematic



Fit result is not very sensitive to J = 1, 2 normalisation as  $\sigma \sim \Delta M_{12}$ 

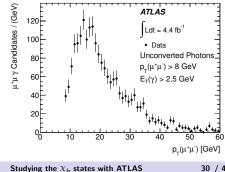
- Bye eye, difficult to notice any difference in the shape of the composite PDF!
- ▶ Reflected in small systematic shift in measured mass (±5 MeV)

#### **Relative Acceptance**



Large difference in acceptance between the three states due to minimum  $p_T(\gamma)$ reconstruction threshold:

- Converted photons have much larger acceptance for all decays at low  $p_T(\Upsilon)$
- Unconverted photons have a much reduced acceptance due to the  $p_T(\gamma) > 2.5$  GeV reconstruction threshold

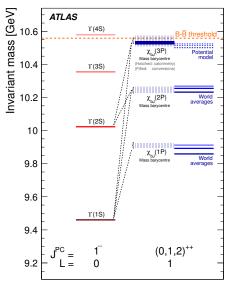


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## Summary

- The known χ<sub>b</sub>(1, 2P) states are observed in radiative decays to Υ(1S) γ
- A new structure at a higher mass is also observed in the Υ(1S) γ and Υ(2S) γ spectra
- The interpretation of this as the \$\chi\_b(3P)\$ states is consistent with theoretical predictions
- The mass of the structure is measured with two separate analyses using converted and unconverted photons with compatible results
- The mass measurement with smaller systematic uncertainties from the converted photon analysis is chosen to represent the final measurement

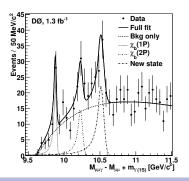
Observed bottomonium radiative decays in ATLAS, L = 4.4 fb<sup>1</sup>



## Confirmation by DØ

Shortly after the publication of the ATLAS result, the  $D\emptyset$  collaboration confirmed the observation of a new structure in the  $\Upsilon(1S)\gamma$  mass spectrum:

Observation of a narrow state decaying into  $\Upsilon(1S) + \gamma$ in  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV Phys. Rev. D 86, 031103(R) (2012) (arXiv:1203.6034 [hep-ex])



"...a third peak is observed at a mass consistent with the new state observed by the ATLAS collaboration."

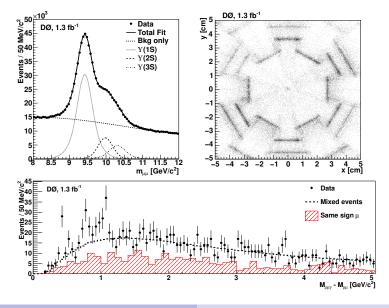
 $\overline{m}_3 = 10.551 \pm 0.014 \text{ (stat.)} \pm 0.017 \text{ (syst.) GeV}$ 

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## Confirmation by DØ

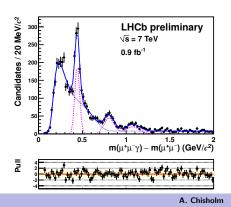


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## Confirmation by LHCb

LHCb later confirmed the observation at ICHEP2012 in a preliminary conference note:

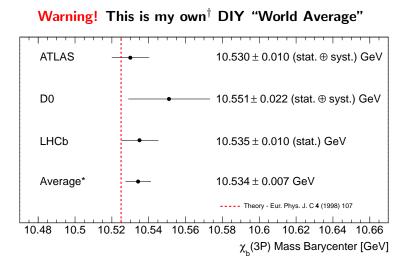
Observation of the  $\chi_b(3P)$  state at LHCb in *pp* collisions at  $\sqrt{s} = 7$  TeV LHCb-CONF-2012-020



"Three peaks are clearly visible, corresponding to the  $\chi_b(1P)$ ,  $\chi_b(2P)$ , and the new  $\chi_b(3P)$  state recently observed by the ATLAS experiment and confirmed by  $D\emptyset$ ."

 $\overline{m}_3 = 10.535 \pm 0.010 \text{ (stat.) GeV}$ 

Studying the  $\chi_b$  states with ATLAS



<sup>†</sup>Nothing to do with the PDG, ATLAS, *DØ* or LHCb!

## Some Renewed Theoretical Interest

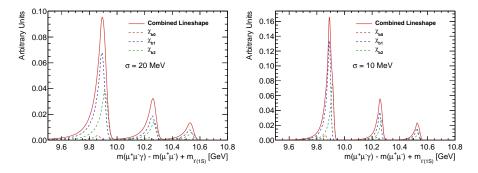
Summary of theoretical work prompted by the observation of the  $\chi_b(3P)$  candidate:

- Potential model results for the newly discovered  $\chi_b(3P)$  states
- arXiv:1201.4096
- Production of X<sub>b</sub> mesons at LHC
- arXiv:1203.4893
- Comment on "Observation of a New  $\chi_b$  State in Radiative Transitions to  $\Upsilon(1S)$  and  $\Upsilon(2S)$  at ATLAS"
- arXiv:1204.1984
- Developments in heavy quarkonium spectroscopy
- arXiv:1205.4189
- $\chi_b(3P)$  splitting predictions in potential models
- arXiv:1208.2186

#### Perhaps the most important implication:

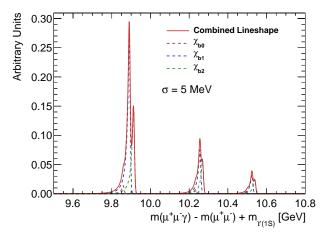
- Another source of feed down into the inclusive  $\Upsilon(nS)$  cross section
- The inclusive Υ(3S) cross section was previously thought to be free from significant feed down, B(X<sub>b</sub>(3P) → Υ(3S)γ) expected to be large

#### You may need to squint!



- Investigate with a crude model and purely qualitative analysis!
- Relative normalisation of J = 0, 1, 2 taken from PDG branching fractions (assume equal production ratio and assume 2P values for 3P)
- ▶ Relative production of n = 1, 2, 3 taken from recent ATLAS  $\Upsilon$  production paper
- $\blacktriangleright$  No acceptance effects, resolution modelled by CB with various different Gaussian widths  $\sigma$





 Mass resolution of around 5 MeV very challenging at the LHC (for ATLAS at least!)

# What next?

# Important to look in other channels!

The (non-)observation of the new state in other decays could shed more light on its nature and confirm / rule out the  $\chi_b(3P)$  interpretation. Some channels that might be possible at the LHC include:

- $\blacktriangleright \ \chi_b \to \Upsilon + \omega$
- $\blacktriangleright \ \chi_b \to \Upsilon + \Phi$
- $\blacktriangleright \ \chi_b \to J/\psi J/\psi$
- Other VV final states?

Cross section (1, 2P) and 3P  $\sigma \cdot \mathcal{B}(\chi_b(3P) \to \Upsilon(1S)\gamma)$  measurements

Possible with 2012 dataset

# Spin, Parity and Polarization measurements

- Likely to require more data, complex analyses...
- ► Polarization is accessible through an angular analysis of final state di-muons in X<sub>b</sub> → Υ(nS)γ (arXiv:1103.4882)

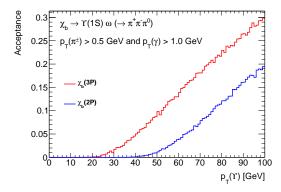
 $\chi_b \to \Upsilon(1S) \, \omega$ 

**CLEO** has observed  $\chi_b(2P) \rightarrow \Upsilon(1S) \omega$  decays (Phys. Rev. Lett. 92, 222002 (2004))

• Only  $\chi_{b1}(2P)$  and  $\chi_{b2}(2P)$  are above the  $\Upsilon(1S)\omega$  threshold

$$\blacktriangleright \ \mathcal{B}(\chi_b(2P) \to \Upsilon(1S)\,\omega) = 1 - 2\% ! \checkmark$$

•  $\omega$  momentum in  $\chi_b$  rest frame only 135(94) MeV! X



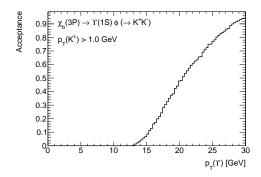
**Bottom Line** - Low acceptance with ATLAS, huge backgrounds! Possible with LHCb?

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# $\chi_b \to \Upsilon(1S) \phi(1020)$

The new state at 10.53 GeV is above the  $\Upsilon(1S) \phi(1020)$  threshold...

- Not yet observed...
- ▶  $\mathcal{B}(\chi_b(3P) \to \Upsilon(1S) \phi)$  not measured or calculated! (as far as I am aware) ×
- ▶ High acceptance with  $\phi \to K^+K^-$  (  $\mathcal{B}(\phi \to K^+K^-)$  is also large ~ 50%) √
- $\Upsilon(1S) + 2$  tracks with ATLAS's limited PID might be messy! X



**Bottom Line** -  $\mathcal{B}$  may be very low! LHCb (with good PID) more sensitive?

 $\chi_{\rm b} \to {\rm J}/\psi \, {\rm J}/\psi$ 

#### $\chi_b \to J/\psi \, J/\psi$ : a potentially a very clean signal at the LHC

- ▶ 4 Lepton (4 $\mu$  more realistic at low  $p_T$ ) very clean, low background  $\checkmark$
- ▶  $\chi_{b1} \rightarrow J/\psi J/\psi$  is Landau-Yang forbidden,  $\Delta m_{0,2} \approx 53$  MeV large enough to be resolved?  $\checkmark$
- ▶  $\mathcal{B}(\chi_{b0} \rightarrow J/\psi J/\psi) \approx 2 \times 10^{-4}$  (Phys. Rev. **D72** (2005) 094018) 🗡

$$\blacktriangleright~ {\cal B}(J/\psi o \mu^+\mu^-)^2 pprox 3.6 imes 10^{-3}$$
 🗡

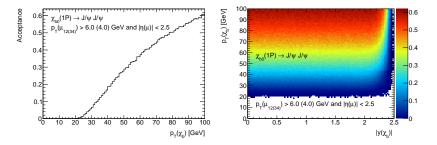
The inclusive cross section for  $\chi_{b0}$  at the LHC is estimated to be as much as  $\sigma(pp \rightarrow \chi_{b0} + X) \approx 1 \mu b$ :

- Rough estimate of 700 events per fb<sup>-1</sup> (before trigger, acceptance and reconstruction)
- Strong potential for observation with > 20fb<sup>-1</sup> from 2011 + 2012 (ATLAS / CMS)

 $\chi_b \to J/\psi J/\psi$ 

Estimates suggest a large raw event yield, but how many  $\chi_b$  are likely to be reconstructed?

▶ Require  $p_T(\mu) > 6$  GeV for triggered  $J/\psi$  and  $p_T(\mu) > 4$  GeV for the other with all muons required to be within  $|\eta(\mu)| < 2.5$ 



- Acceptance and trigger thresholds allow only very boosted ( $p_T > 20 GeV$ )  $\chi_b$  to be reconstructed
- $\blacktriangleright$  This will significantly (factor of  $\sim$  100) reduce the yields on the previous slide!

**Bottom Line** - Likely to need 10s fb<sup>-1</sup> data for an observation, possible in 2012?

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|-------------|---|---------|
|-------------|---|---------|

## Conclusion

- The known X<sub>b</sub>(1, 2P) states are observed in radiative decays to Υ(1S) γ at ATLAS
- A new structure at a higher mass is also observed in the Υ(1S) γ and Υ(2S) γ spectra
- ► The interpretation of this as the X<sub>b</sub>(3P) states is consistent with theoretical predictions
- Many more interesting opportunities at the LHC!

## Thank you for listening!

 $\chi_{\textit{c}}$  at the LHC

