The W mass measurement at LHCb

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Outline

Introduction to EW physics

Constraining EW observables

Overview of the current measurements and experiments

The W mass measurement at LHCb

EW physics at LHCb

The analysis of 2016 data

Ongoing studies with the full Run 2 dataset

Prospects for the future

Electroweak theory

The Electroweak theory

Main magnitudes ruling EW interactions are related to each other:

Abdus Salam, Steven Weinberg and Sheldon Lee Glashow

$$
m_W^2 \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2}G_F}(1 + \Delta)
$$

$$
\Gamma_W \propto G_F m_W^3
$$
 Higher order corrections

The global EW fit

Global fits to EW observables allow to test current (and new) theoretical model(s)

Past and present of the W mass measurement

7

A wider theoretical picture

- Fundamental magnitude related to other EW observables
- The experimental sensitivity is still away from the theoretical best fit 12 MeV / 7 MeV
- Interesting implications in BSM models with other magnitudes of interest

Total uncertainty Stat. uncertainty

ALEPH DELPHI $L₃$ **OPAL**

[\[JHEP 01 \(2022\) 036\]](https://doi.org/10.1007/JHEP01(2022)036)

The W mass measurement at LHCb

Production mechanism

- A proton-proton collider is more challenging to measure the W mass:
	- W bosons are produced in a mixture of positive and negative helicity states
	- Must accurately describe the angular cross-section (larger uncertainties)
	- More backgrounds through heavy-flavour processes
- But much higher total production cross-section and larger calibration samples
	- One of the main objectives is being able to extrapolate the *Z* measurements to the *W*.

Related detector features

- Detector in the forward region with excellent momentum and vertex resolutions
- Coverage is complementary to ATLAS and CMS (with some overlapping at low pseudorapidity)

W and Z production at LHCb Example 2 production at LHCb

- Z decays constitute the most natural way of controlling muons from W decays and the cross-section
- Anti-correlation of the PDF uncertainties at low Bjorken-x allows achieving a similar precision of the LHC experiments to the theoretical best fit for the W mass [Eur. Phys. J. C 75, 601 \(2015\)](https://arxiv.org/abs/1508.06954v2)

Anti-correlation of uncertainties from PDFs

[Eur. Phys. J. C 75, 601 \(2015\)](https://arxiv.org/abs/1508.06954v2)

Single event signature

Results from other experiments

 $m_W = 80387 \pm 12_{\rm stat} \pm 15_{\rm syst} \rm MeV$

- Barrel-like detectors allow to measure missing transverse energy and the transverse mass
	- Measurement can be done measuring different quantities
- In modern experiments, a similar sensitivity can be obtained measuring the momentum of the outgoing lepton

Analysis strategy

- Carefully measure the muon transverse momentum
- Use plain LHCb Pythia8 simulation and reweight using samples with generator-level information from different models
- Corrections due to the efficiencies of the different selection steps (reconstruction, trigger, topological, offline selection)
- Study and determine background from simulation (except for the contribution from hadrons originating decays-in-flight)

Detector alignment and calibration

- The LHCb trigger changed significantly for Run 2
- Real-time alignment and calibration can be optimized offline for EW studies
- Need to re-process the data using dedicated tools
- Apply corrections and smearing to simulation to account for subtle effects that significantly affect the momenta distributions

Calibration using muons

Charge-dependent curvature biases

- The analysis relies highly on the detector alignment
	- Misalignment of 10µm translates into a O(50MeV) shift
- Default LHCb alignment and calibration not suitable to study candidates with high transverse momentum
- Need to re-run the alignment and calibration offline using Z
- Avoid double bias from the momentum resolution using the pseudo-mass method

$$
M^{\pm} = \sqrt{2p^{\pm}p_T^{\pm}\frac{p^{\mp}}{p_T^{\mp}}}(1-\cos\theta)_{\text{In}}
$$

spired by [Phys. Rev. D 91, 072002](https://journals.aps.org/prd/abstract/10.1103/PhysRevD.91.072002)

Miguel Ramos Pernas The W mass measurement at LHCb 16/03/2022 19

Charge-dependent curvature biases

Fit the asymmetries to the pseudomass and translate this into shifts in q/p

[\[JHEP 01 \(2022\) 036\]](https://doi.org/10.1007/JHEP01(2022)036), [\[LHCB-PAPER-2021-024\]](https://cds.cern.ch/record/2780004)

Corrections to the simulation

slide)

The W cross-section

$$
\frac{d\sigma}{dp_T^W dy dM d \cos \vartheta d\varphi} = \frac{3}{16\pi} \frac{d\sigma^{unpol.}}{dp_T^W dy dM}
$$
\n
$$
(At order α_s^2)\n
$$
+ A_2 \frac{1}{2} \sin^2 \vartheta \cos 2\varphi + A_3 \sin \vartheta \cos \varphi + A_4 \cos \vartheta + A_5 \sin^2 \vartheta \sin 2\varphi + A_6 \sin 2\vartheta \sin \varphi + A_7 \sin \vartheta \sin \varphi \}
$$
\n
$$
+ A_5 \sin^2 \vartheta \sin 2\varphi + A_6 \sin 2\vartheta \sin \varphi + A_7 \sin \vartheta \sin \varphi \}
$$
$$

Simulating signal decays

[\[JHEP 01 \(2022\) 036\]](https://doi.org/10.1007/JHEP01(2022)036), [\[LHCB-PAPER-2021-024\]](https://cds.cern.ch/record/2780004)

- POWHEG + Pythia gives the best description of the unpolarized cross-section and is chosen as the baseline generator
	- Varied success with other generators, used to determine systematic uncertainties
- DYTurbo performs well at reproducing the angular cross-section

Modelling the W boson transverse momentum

The limited knowledge on the transverse momentum of the W bosons can be compensated by floating QCD floating parameters [\[arXiv:1907.09958\]](https://arxiv.org/abs/1907.09958)

Modelling the boson transverse momentum

[\[arXiv:2112.07458 \(submitted to](https://arxiv.org/abs/2112.07458) JHEP)]

 $\frac{d\sigma}{dp_{_{\rm T}}^2} \, [\mathrm{pb}/(\mathrm{GeV}/c)]$ • The momentum of the outgoing 10 **BREAD** muon is strictly related to that of the boson **LHCb** 5.1 fb^{-1} Must ensure the correlation is $\sqrt{s} = 13 \text{ TeV}$ Φ 0 maintained after the fit ○ Fit *Z* variables simultaneously **Statistical Uncertainty Total Uncertainty Resbos** 10^{-1} **Pythia, LHCb tune** $\phi^* \equiv \arctan \left(\frac{\pi - \Delta \phi}{2} \right) / \cosh \left(\frac{\Delta \eta}{2} \right) \sim \frac{p_T}{M}$ **POWHEG+Pythia MatchBox** $10²$ 10 [\[EPJC 71, 1600 \(2011\)\]](https://arxiv.org/abs/1009.1580) $p_{\rm T}^{\rm Z}$ [GeV/c]

Polarized cross-section

- Angular part is better described with DYTurbo
- However, there angular coefficients suffer low accuracy at low transverse momentum values [\[JHEP 11 \(2017\) 003\]](https://arxiv.org/abs/1708.00008)
- Uncertainties from DYTurbo mitigated by floating A_3
	- Otherwise the uncertainty would be O(20 MeV)

The simulation process (PDF set)

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Selections

- EW physics with leptons in the final state can be done at LHCb with simple selections based on the transverse momentum, impact parameter, isolation and particle identification
- Selection biases studied in data and simulation for Z and Υ(1S) decays (isolation biases only studied in the former)
	- Associated systematic uncertainties determined by varying the binning scheme, parametrizations and selections

$$
\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2 \big(\mathrm{rad}^{-2} \big)}
$$

Determining the efficiencies

Three main sources of acceptance biases:

● Trigger efficiencies

Trigger efficiency

 $i/2$ $\overline{}$

 $\overline{+}$

 $= 10, 4.40 < \eta < 5.00$ $= 9$ 4.12 < $n \le 4.40$ $i = 8, 3.85 < \eta < 4.12$

 $i = 7, 3.58 < \eta < 3.85$

 $i = 6, 3.30 < \eta < 3.58$

 $-\pi < \phi < -\pi/2$

LHCb 1.7 fb $^{-1}$

Backgrounds

- Most of them modelled from dedicated simulated samples
	- \circ Single-top, quark/anti-quark (t, b, c), Z/W decays, Drell-Yan
	- Cross-sections normalized to the W
- Description of the QCD background (decays-in-flight) obtained from data
	- Sample with inverted muon-identification requirements
	- Weight and parametrize the data using a Hagedorn distribution
- Accurately describes the Jacobian peak (region with highest sensitivity to $m_{_W}\!)$

Miguel Ramos Pernas The W mass measurement at LHCb 16/03/2022 16/03/2022 30

Systematic uncertainties

Fit to extract the W mass FILE TO THEP 01 (2022) 036], [\[LHCB-PAPER-2021-024\]](https://cds.cern.ch/record/2780004)

- 5D-weighted fit using the Beeston-Barlow approach
- Fit simultaneously *W* and *Z* data

LHCb measures the W mass!

- Measurement of the W mass using 2016 data
- Published on January 2022
- Shows the LHCb capabilities of doing high-precision measurements

Prospects for the future

What can we do in the near future?

Is including 2017 and 2018 data straight-forward?

- It is straight-forward, but we must ask ourselves the following questions:
	- Can we optimize any part of the analysis strategy?
	- Can we use any of the new options available in the market?
	- Are there ways to make the result more accessible/easy to use for people outside the collaboration?
- The result using 2016 data shows the capabilities of the LHCb detector to contribute to this measurement, but it is worth re-considering our strategy before studying the full Run 2 data sample

Improving the simulation

- Take advantage of the latest developments on the theory side
	- Switch to more accurate predictors of the boson production
	- New PDF sets (NNPDF 4.0)
- Change the treatment generators / PDF sets when calculating systematic uncertainties
	- Drop known inaccurate PDF sets
	- Revisit the way to handle the different predictors and the order of the accuracy (NLL, NNLL, …)
- Ongoing studies, feedback is really welcome!

Towards doing an unfolded measurement

[\[JHEP 01 \(2022\) 036\]](https://doi.org/10.1007/JHEP01(2022)036), [\[LHCB-PAPER-2021-024\]](https://cds.cern.ch/record/2780004)

- Ongoing studies to see if we can publish the unfolded transverse momentum distribution
- Facilitate comparing prediction and observables
- Quite challenging from the experimental point of view:
	- Must have a good control of the backgrounds (especially in the selection variables)
	- The systematic uncertainties might turn much bigger with the unfolding methods

Expected sensitivity for the full Run 2 analysis

- We expect to reduce the overall experimental uncertainty to 15 MeV
- The analysis becomes systematically dominated
	- A more careful description of the physics is necessary
- Eager to see the result of combining the measurements of all the LHC experiments

 $m_W = 80354 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 9_{\text{PDF}}$ MeV

Summary

Summary

$m_W = 80354 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 9_{\text{PDF}}$ MeV

Looking forward to hearing your comments and suggestions

The W mass measurement using 2016 data is a big milestone at LHCb

- Already exploring new strategies to improve the result with the full Run 2 data sample
- Improvements on the physics modelling are strictly necessary to be competitive

Thank you!